

PARTICLE DYNAMICS IN LOW-ENERGY TRAVELLING-WAVE LINACS

J.M. Corstens, A.F.J. Hammen, J.I.M. Botman

Eindhoven University of Technology, Cyclotron Laboratory,
P.O.Box 513, 5600 MB Eindhoven, The Netherlands.

Abstract

In this paper, we present the particle dynamics in low-energy travelling-wave linear accelerators, applying analytical theory, based on Hamiltonian mechanics, and numerical simulations, performed by commercially available codes. The paper is an extension on earlier work, presented at EPAC'98. Cylindrical co-ordinates are used and solenoid magnetic fields are incorporated. The Hamiltonian equations of motion are given and examples of calculations are presented and compared to numerical simulations, yielding excellent agreement between both approaches.

1 INTRODUCTION

At the Eindhoven University of Technology, a Hamiltonian theory of particle motion of accelerated electrons in standing-wave RF structures has been developed[1]. However, this theory was only valid in principle under the assumption that the particle velocity equals the phase velocity of the main accelerating wave. This condition is not fulfilled in low-energy linear accelerators. The present, adapted Hamiltonian theory, in which also solenoidal magnetic fields are incorporated does not have this restriction. The set of equations obtained with this Hamiltonian theory is suitable for particle motion calculation. Calculations using this set of equations are compared to the results of a commercially available particle tracking code, yielding excellent agreement.

2 THE EQUATIONS OF MOTION

The vector potential in cylindrical co-ordinates representing electromagnetic waves in the periodic part of the linac is given by [2]:

$$\begin{aligned} A_r &= \sum_n \frac{k_n}{\omega} \frac{a_n E_z(z)}{\alpha_n} I_1(\alpha_n r) \cos(k_n z - \omega t), \\ A_\theta &= \sum_{j=1}^8 A_{\theta,j}, \\ A_z &= -\sum_n \frac{a_n E_z(z)}{\omega} I_0(\alpha_n r) \sin(k_n z - \omega t), \end{aligned} \quad (1)$$

in which \sum_n stands for $\sum_{n=-\infty}^{\infty}$, a_n for the Fourier coefficients, $k_n = k_f + 2\pi n/d$ with d the cell length and k_f the phase shift per cell, $\alpha_n^2 = k_n^2 - k^2$, $E_z(z)$ the amplitude of the electric field in the z -direction. Note that this amplitude is a

function of the longitudinal co-ordinate z . $A_{\theta,j}$ is the vector potential of the static magnetic field j . This vector potential of solenoid j , expanded up to third order in r , reads:

$$A_{\theta,j}(r, z) = \frac{1}{2} r B_z(r=0, z) - \frac{1}{16} r^3 \frac{\partial^2}{\partial z^2} (B_z(r=0, z)), \quad (2)$$

$$\text{with: } B_z(r=0, z) = -\mu_0 \frac{\partial}{\partial z} V_m(z),$$

in which:

$$V_m(z) = -\frac{NI}{2L} \left(\sqrt{\left(z + \frac{L}{2}\right)^2 + a^2} - \sqrt{\left(z - \frac{L}{2}\right)^2 + a^2} \right),$$

with μ_0 the magnetic permeability in vacuum, I the current through the solenoid, N the number of turns, L the length and a the radius of the solenoid. z is defined at the axis with respect to the center of the solenoid.

Relativistic motion is described by the Hamiltonian:

$$\begin{aligned} H &= \{E_r^2 + (p_r - eA_r)^2 c^2 + \left(\frac{p_\theta}{r} - eA_\theta\right)^2 c^2 + \\ &\quad (p_z - eA_z)^2 c^2\}^{1/2}, \end{aligned} \quad (3)$$

where c is the speed of light, E_r the particle rest energy, p_r, p_θ, p_z the canonic momenta in cylindrical co-ordinates and A_r, A_θ, A_z the components in cylindrical co-ordinates of the vectorpotential, given by eq. 1. The new Hamiltonian becomes $K_1 = -p_z, -H$ and t form a new pair of conjugated canonical variables. On this Hamiltonian K_1 the following scaling transformations is performed:

$$\begin{aligned} \pi_r &= \frac{c}{H_i} p_r, \pi_\theta = \frac{c}{H_i} p_\theta, \pi_z = \frac{c}{H_i} p_z, h = -\frac{H}{H_i}, \\ e_r &= \frac{E_r}{H_i}, \epsilon_n = \frac{e a_n E_z(z)}{H_i k}, \varsigma = ct. \end{aligned} \quad (4)$$

in which H_i is the initial energy, yielding Hamiltonian $K_2 = -\pi_z$:

$$K_2 = -\left\{ h^2 - e_r^2 - (\pi_r - eA_r^*)^2 - \left(\frac{\pi_\theta}{r} - eA_\theta^*\right)^2 \right\}^{1/2} - eA_z^*, \quad (5)$$

in which: $\vec{A}^* = \frac{c}{H_i} \vec{A}$.

In eq. 5 ζ is transformed to $\zeta_2 = \zeta - k_f z/k$ (new variables are indicated by higher numerical subscripts), by generating function $G_0 = -h\zeta_2 - h k_f z/k$ ($h_2 = h$). Further the signs of h_2 and ζ_2 are changed in h_3 and ζ_3 by generating function

$G_1 = h_2 \zeta_3$. Now h_3 is a positive quantity and $k\zeta_3$ has the meaning of phase. The final Hamiltonian K_f becomes (in new co-ordinates, numerical subscripts are omitted):

$$K_f = K_2 + \frac{\partial G_0}{\partial z} + \frac{\partial G_1}{\partial z} = \frac{k_f}{k} h - \{h^2 - e_r^2 - (\pi_r - eA_r^*)^2 - \left(\frac{\pi_\theta}{r} - eA_\theta^*\right)^2\}^{1/2} - eA_z^*, \quad (6)$$

in which the arguments of the components of the vector potential ($k_n z - \omega t$) in eq. 1 has been replaced by $(2\pi n z / d + k\zeta)$. The equations of scaled energy and phase are:

$$\begin{aligned} \frac{dh}{dz} &= -\frac{\partial K_f}{\partial \zeta} = \frac{\pi_{r,kin}}{\pi_{z,kin}} \left(-\sum_n \frac{\varepsilon_n}{\alpha_n} k_n k I_1(\alpha_n r) \sin\left(\frac{2\pi n}{d} z + k\zeta\right) \right) - \\ &\quad \sum_n \varepsilon_n k I_0(\alpha_n r) \cos\left(\frac{2\pi n}{d} z + k\zeta\right), \\ \frac{d\zeta}{dz} &= \frac{\partial K_f}{\partial h} = -\frac{h}{\pi_{z,kin}} + \frac{k_f}{k}, \end{aligned} \quad (7)$$

in which:

$$\pi_{z,kin} = \sqrt{h^2 - e_r^2 - \pi_{r,kin}^2 - \left(\frac{\pi_{\theta,kin}}{r}\right)^2}.$$

The transversal equation of motion for the canonical momentum reads:

$$\begin{aligned} \frac{d\pi_{r,can}}{dz} &= -\frac{\partial K_1}{\partial r} = \frac{\pi_{r,kin}}{\pi_{z,kin}} e \frac{\partial A_r^*}{\partial r} + e \frac{\partial A_z^*}{\partial r} + \frac{\pi_{\theta,kin}^2}{r^3 \pi_{z,kin}} + \\ &\quad e \frac{\pi_{\theta,kin}}{r^2 \pi_{z,kin}} \frac{\partial(rA_\theta^*)}{\partial r}. \end{aligned} \quad (8)$$

This expression is inconvenient for practical use. Furthermore, $\pi_{r,can}$ is not of interest but $\pi_{r,kin}$, therefore the differential equation for $\pi_{r,kin}$ will be derived in the appendix and is given by:

$$\begin{aligned} \frac{d\pi_{r,kin}}{dz} &= \sum_n \varepsilon_n \frac{I_1(\alpha_n r)}{\alpha_n} k \left(k - k_n \frac{h}{\pi_{z,kin}} \right) \sin\left(\frac{2\pi n}{d} z + k\zeta\right) + \\ &\quad \frac{\pi_{\theta,kin}^2}{r^3 \pi_{z,kin}} + e \frac{\pi_{\theta,kin}}{r^2 \pi_{z,kin}} \frac{\partial(rA_\theta^*)}{\partial r}. \end{aligned} \quad (9)$$

The remaining transversal equations are:

$$\begin{aligned} \frac{dr}{dz} &= \frac{\partial K_f}{\partial \pi_{r,can}} = \frac{\pi_{r,kin}}{\pi_{z,kin}}, \\ \pi_{\theta,kin} &= \pi_{\theta,can}(0) - e r A_\theta^*(r, z), \\ \frac{d\theta}{dz} &= \frac{\partial K_f}{\partial \pi_\theta} = \frac{\pi_{\theta,kin}}{r^2 \pi_{z,kin}}. \end{aligned} \quad (10)$$

It is seen that $\pi_{r,kin}$ is needed in the set of equations and not $\pi_{r,can}$. Furthermore, the expression $d\pi_{r,kin}/dz$ is much simpler than the expression $d\pi_{r,can}/dz$. With this set of equations particle motion calculations can be performed.

The equations of motion for calculations without static external solenoidal magnetic field are obtained by setting A_θ equal to zero. Calculations without electromagnetic field are obtained by setting $\varepsilon_n(z)$ equal to zero. In this case the differential equation for energy changes in $h(z)=1$ and the differential equation for the phase does not have any meaning anymore, because of the absence of an electromagnetic field.

3 CALCULATIONS COMPARED TO SIMULATIONS

In this section calculations based on the Hamiltonian equations of motion of the previous section are presented and compared to the results of the particle-tracking code General Particle Tracer (GPT) [3]. In this code, the same electric and magnetic field description is specified as is used in the Hamiltonian calculations, however this particle-tracking code uses an entirely different calculation method. Calculations of motion in an electromagnetic field with and without static external magnetic field have been applied to the periodic part of the Eindhoven 10 MeV linac.

The energy and phase have been calculated for the 10 MeV linac using eq. 7 (fig. 1). Because of the small dependence of the particle distance to the z -axis on the calculation of energy and phase, particles with a different initial position and divergence will have similar energy and phase evolution. Calculations are performed at four different injection phases, with an initial energy of 1 MeV. Energy calculations are compared to the results of GPT, showing good agreement between both methods. The particle phase is not provided by GPT, so a comparison of phases is not possible. However, because of the agreement in energy calculations, it is obvious that the Hamiltonian phase calculations must be consistent.

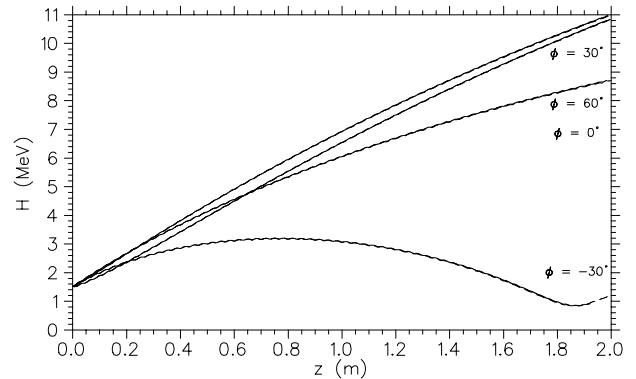


Figure 1: Particle energy H as a function of longitudinal co-ordinate z in the periodic part of the linac, at various injection phases. Hamilton calculations are presented by the dashed line, GPT simulations by the solid line.

Fig. 2 depicts the calculations of the distance of a particle to the z -axis, r , and the divergence r' in the linac, using eqs. 9 and 10. Particles are injected at various phases at 1 MeV parallel to the z -axis with an initial radius $r = 3\text{mm}$. It is shown that there is a good agreement between Hamiltonian calculations and GPT simulations.

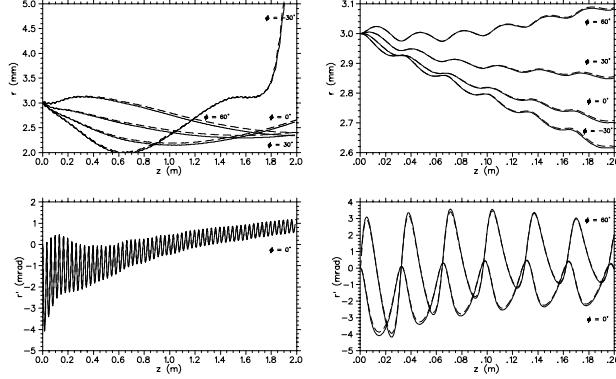


Figure 2: Particle distance to z -axis, r , and divergence, r' , calculations as a function of z , for the entire periodic part of the linac (left figs.) and for the first six cells of the periodic part (right figs.) Hamilton calculations are presented by the dashed line, GPT simulations by the solid line.

Fig. 3 shows the azimuthal co-ordinate θ and its derivative θ' as a function of z in the linac for the same calculations as depicted in fig. 2. Again there is a good agreement between Hamiltonian calculations and GPT simulations.

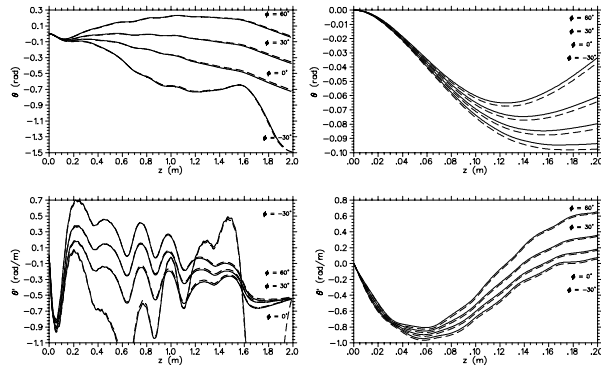


Figure 3: Particle azimuthal co-ordinate, θ , and its derivative, θ' , calculations as a function of z , for the entire periodic part of the linac (left figs.) and for the first six cells of the periodic part (right figs.) Hamilton calculations are presented by the dashed line, GPT simulations by the solid line.

4 CONCLUSIONS

An analytic description of particle motion in linear accelerators has been developed, in an earlier paper restricted to higher energy standing wave structures,

while in the present paper this has been extended to the case of low-energy travelling-wave linacs, in which solenoidal fields may be incorporated. The Hamiltonian calculations yield a remarkable agreement with results of particle-tracking codes but they are much faster performed than simulations by these codes.

5 REFERENCES

- [1] A.F.J. Hammen, J.M. Corstens, J.I.M. Botman, H.L. Hagedoorn, W.H.C. Theuws, "Hamiltonian Calculations on Particle Motion in Linear Electron Accelerators", European Particle Accelerator Conference, Stockholm, Sweden (1998).
- [2] R.W. de Leeuw, "The accelerator injection chain of the electron storage ring EUTERPE", Ph.D. Thesis, Eindhoven University of Technology, Eindhoven, The Netherlands (1996).
- [3] General Particle Tracer, "User Manual", Pulsar Physics, Utrecht, The Netherlands (1997).

APPENDIX

Consider the following equations:

$$\begin{aligned} \frac{d\pi_{r,kin}}{dz} &= \frac{d\pi_{r,can}}{dz} - e \frac{dA_r^*}{dz}, \\ \frac{dA_r^*}{dz} &= \frac{\partial A_r^*}{\partial r} \frac{dr}{dz} + \frac{\partial A_r^*}{\partial z} + \frac{\partial A_r^*}{\partial t} \frac{1}{v_z}, \\ \frac{dr}{dz} &= \frac{\pi_{r,kin}}{\pi_{z,kin}}. \end{aligned} \quad (A1)$$

Thus, using eq. 8:

$$\begin{aligned} \frac{d\pi_{r,kin}}{dz} &= e \frac{\partial A_z^*}{\partial r} - e \frac{\partial A_r^*}{\partial z} - e \frac{\partial A_r^*}{\partial t} \frac{1}{v_z} + \frac{\pi_{\theta,kin}^2}{r^3 \pi_{z,kin}} + \\ &e \frac{\pi_{\theta,kin}}{r^2 \pi_{z,kin}} \frac{\partial(rA_\theta^*)}{\partial r}. \end{aligned} \quad (A2)$$

Substitution of the vectorpotential given by eq. 1 yields:

$$\begin{aligned} \frac{d\pi_{r,kin}}{dz} &= \sum_n \epsilon_n I_1(\alpha_n r) \left\{ -\alpha_n + \frac{k_n^2}{\alpha_n} - \frac{\omega}{v_z} \frac{k_n}{\alpha_n} \right\} \sin(k_n z - \omega t) + \\ &\frac{\pi_{\theta,kin}^2}{r^3 \pi_{z,kin}} + e \frac{\pi_{\theta,kin}}{r^2 \pi_{z,kin}} \frac{\partial(rA_\theta^*)}{\partial r}. \end{aligned} \quad (A3)$$

in which the z -dependence of the amplitude in A_r has been ignored. After some mathematical manipulation and replacement of the argument $(k_n z - \omega t)$ by $(2\pi n z/d + k\zeta)$ the following expression holds:

$$\begin{aligned} \frac{d\pi_{r,kin}}{dz} &= \sum_n \epsilon_n \frac{I_1(\alpha_n r)}{\alpha_n} k \left(k - k_n \frac{h}{\pi_{z,kin}} \right) \sin\left(\frac{2\pi n}{d} z + k\zeta\right) + \\ &\frac{\pi_{\theta,kin}^2}{r^3 \pi_{z,kin}} + e \frac{\pi_{\theta,kin}}{r^2 \pi_{z,kin}} \frac{\partial(rA_\theta^*)}{\partial r}. \end{aligned} \quad (A4)$$