

Emittance budget in the transition regime between linear emission and space charge dominated photoemission

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Contents

- Why high-brightness electron injectors for FELs?
- How transition regime of emission (TRE) defined?
- Why TRE interesting?
- How to model dynamics in TRE?
- How the TRE dynamics impacting downstream beam qualities?

FEL-based X-ray facilities require high-brightness electron injectors

- High peak current (I) & low emittance $(\varepsilon_n) \rightarrow$ high beam brightness (B_n)
 - \rightarrow High $I \rightarrow$ high charge and short length \rightarrow high FEL gain and efficiency
 - \rightarrow Low $\varepsilon_n \rightarrow$ required beam energy at a given wavelength (λ)
- Fixed charge → emittance minimization
- Emittance can only be improved in the injector
- Emittance budget & optimization strategy
 - → Minimizing space charge contribution
 - → Improving cathode intrinsic emittance
 - → Making other items negligible
- Intrinsic emittance → lower limit of final emittance

$$B_n \propto \frac{I}{{\varepsilon_n}^2}$$

$$\frac{\varepsilon_n}{\beta\gamma}\approx \frac{\lambda}{4\pi}$$

$$arepsilon_n \propto \sqrt{arepsilon_{th}^2 + arepsilon_{spch}^2 + arepsilon_{rf}^2 + arepsilon_{Bz}^2 + \cdots} + coupling items$$
intrinsic emittance $(arepsilon_{th})$
space charge emittance $(arepsilon_{spch})$
rf emittance $(arepsilon_{rf})$
cathode magnetic field caused emittance $(arepsilon_{Bz})$



W. Decking, H. Weise, Commissioning of the European XFEL accelerator, Paper MOXAA1, IPAC 2017

F. Stephan, M. Krasilnikov, High Brightness Photo Injectors for Brilliant Light Sources, Chap. Of "Synchrotron Light Sources and Free-Electron Lasers", 2016

Ch.-X. Tang, Paper MO2A04, LINAC 2016

F. Sannibale, W.S. on High Repetition-rate XFEL Physics and Technology, 2017

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Emittance optimization at PITZ for FLASH and European XFEL



- Photo Injector Test facility at DESY in Zeuthen (PITZ)
- Typical optimization scheme at PITZ
 - Slit-Scanning emittance vs. gun solenoid current at a given transverse cathode laser spot size
 - Optimize the spot size for smallest achievable emittance

(fixed bunch charge, cathode laser pulse length and shape, gun and booster gradient and phase)

RF Gun¹⁻²

- **L-band** (1.3 GHz) 1.6-cell copper cavity
- Ecath ≥ 60 MV/m → 7 MeV/c e-beams
- ■650 µs ×10 Hz → up to **45 kW** av. RF power
- **•Cs**₂**Te** PC³ (QE~5-20%) \rightarrow up to 6 nC / bunch

¹Phys. Rev. ST Accel. Beams 13, 020704 (2010)

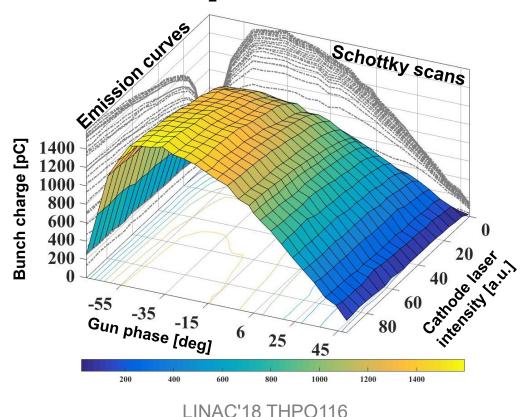
²Phys. Rev. ST Accel. Beams 15, 100701 (2012)

³Cathode production: S. Lederer, L. Monaco, D. Sertore, P. Michelato

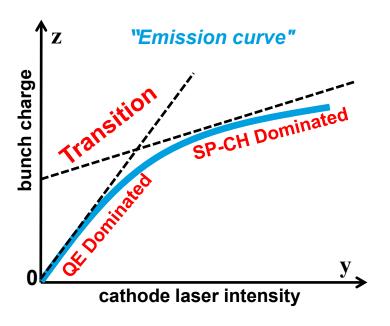
Cathode R&D → next talk of Laura Monaco

Transition regime of photoemission in RF gun environment

Emission characterization in the gun Cs₂Te, 60 MV/m



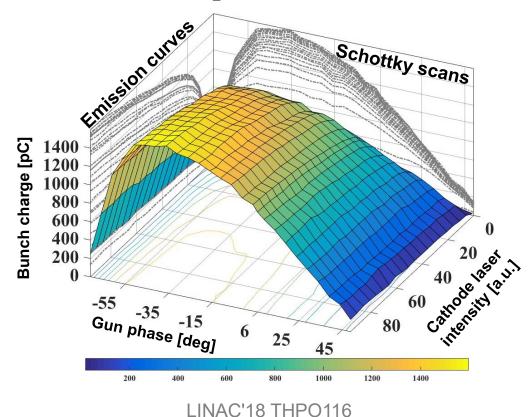
Concept: transition regime



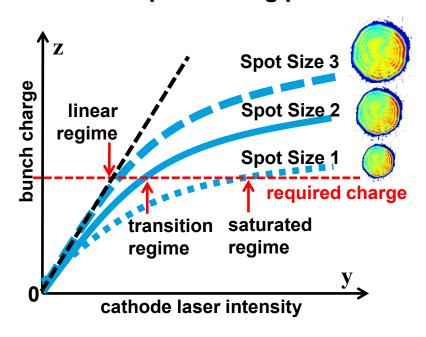
QE: Quantum Efficiency **SP-CH:** Space-Charge

Transition regime of photoemission in RF gun environment

Emission characterization in the gun Cs₂Te, 60 MV/m



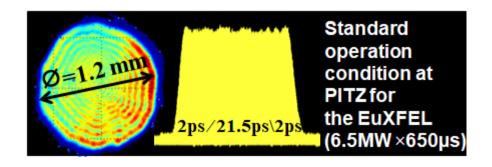
Concept: working point



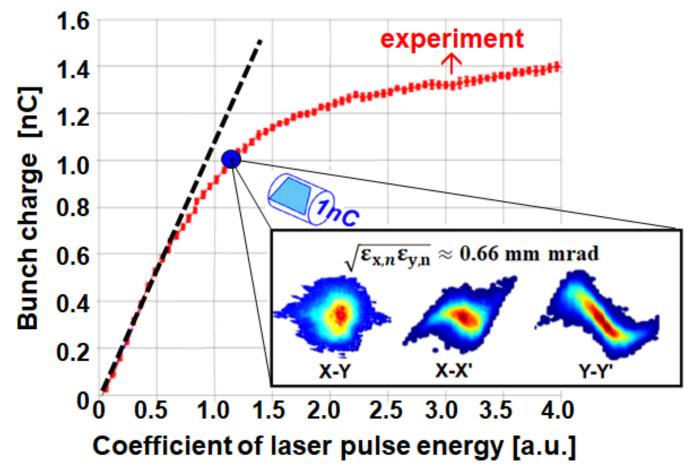
QE: Quantum Efficiency **SP-CH:** Space-Charge

Spot Size: transverse laser spot size on cathode Trans. distributions used only for illustration purpose

Experimental observation on emittance in transition regime of emission

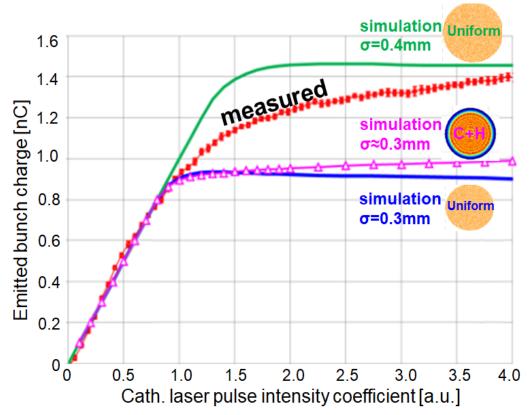


Under standard operation conditions at PITZ, best emittance obtained in transition regime of emission!



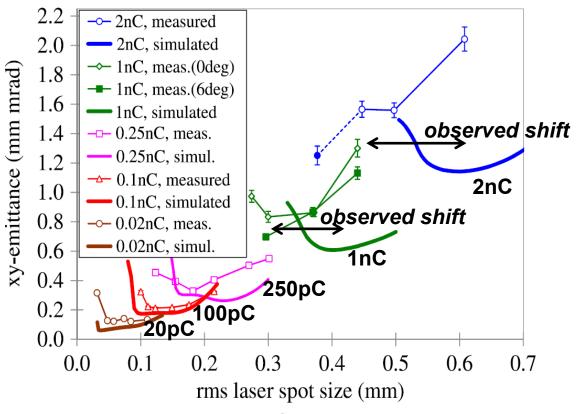
Dynamics in TRE cannot be well reproduced by simulations

Simulated emission curve ≠ measured one



NIM A 889, 129-137 (2018) NIM A 871, 97–104 (2017)

Simulated optimum laser spot size ≠ measured one



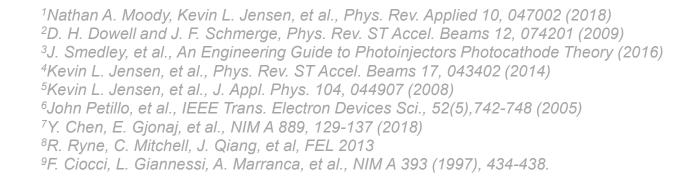
M. Krasilnikov, et al., Phys. Rev. ST Accel. Beams 15, 100701 (2012)

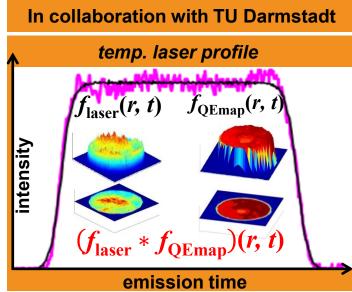
Bring cathode and electron emission physics to beam dynamics

- Not yet straightforward consideration of cathode effects¹-³ in particle simulations
- Emission model needed for particle dynamics with collective effects at cathode
- → first priority: model emission dynamics in strong fields
 - incorporating PIC-ready emission model⁴⁻⁶ with a Lienard-Wiechert approach⁷⁻⁹
 - transient charge packet creation by interplays of cathode QE with time and space dependent rf and beam self-fields

Features

- → measurement-based model training
- → dynamic beam production through cathode physics model
- → taking into account impacts of cathode field effects onto intrinsic beam slice formation





$$dQ(r_{\perp},t) = \frac{e\alpha dE_{las}(r_{\perp},t)dr_{\perp}dt}{\hbar\omega \left\{1 + E_{a}/[\hbar\omega - \Phi_{\text{eff}}(r_{\perp},t)]\right\}^{2}}$$

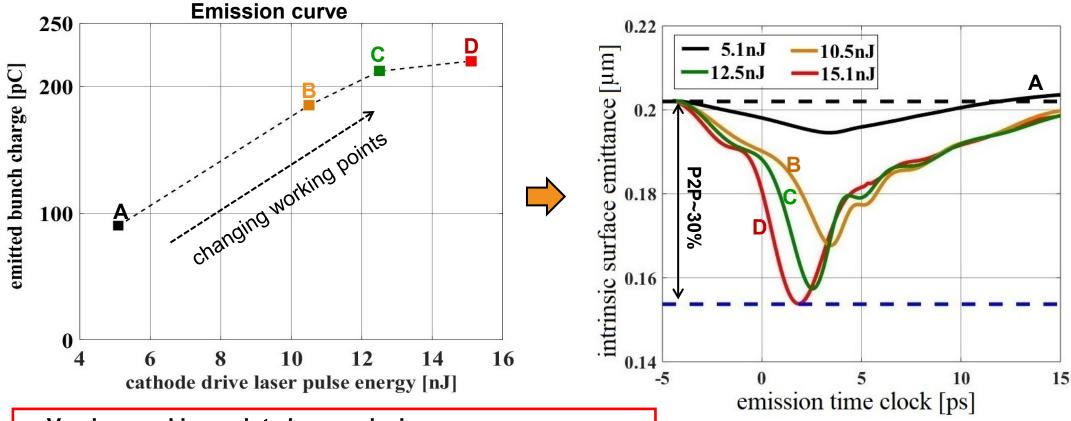
Charge production per time step

Field-dependent cathode work function

$$\Phi_{\text{eff}} = \Phi_0 \pm$$

$$\sqrt{e[E_{\rm rf}(r_\perp,t,z=0)+E_{\rm sc}(r_\perp,t,z=0)]/4\pi\epsilon_0}$$

Effect on intrinsic surface emittance



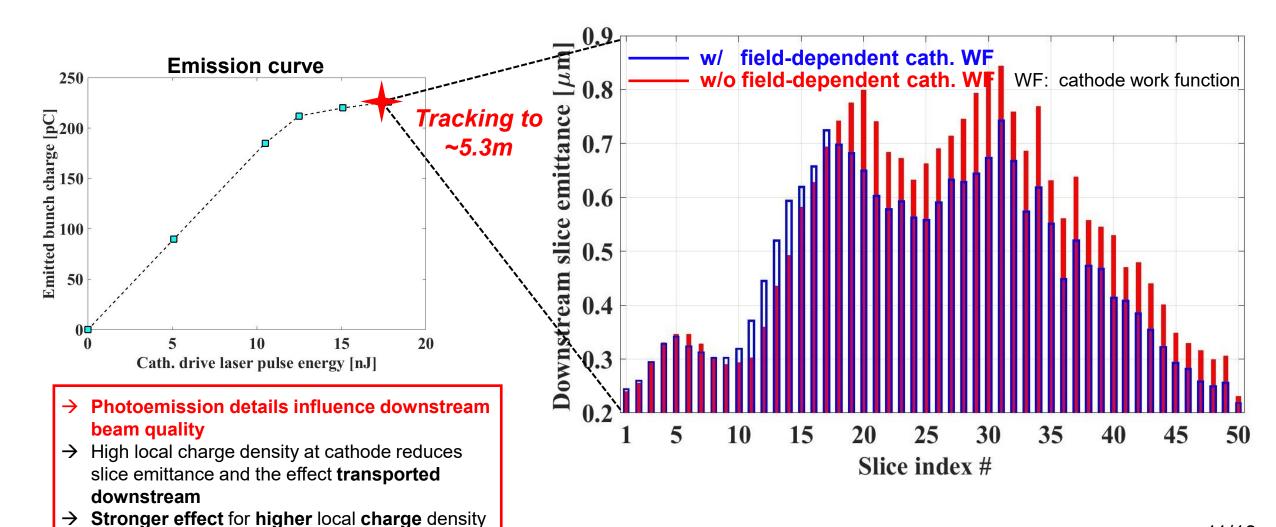
Varying working point along emission curve

- → changing intra-bunch modulation of intrinsic surface emittance
- → overall surface emittance reduction by space charge fields
- → peak to peak ~30% and ~10% in average
- → stronger effects for higher local charge densities at cathode

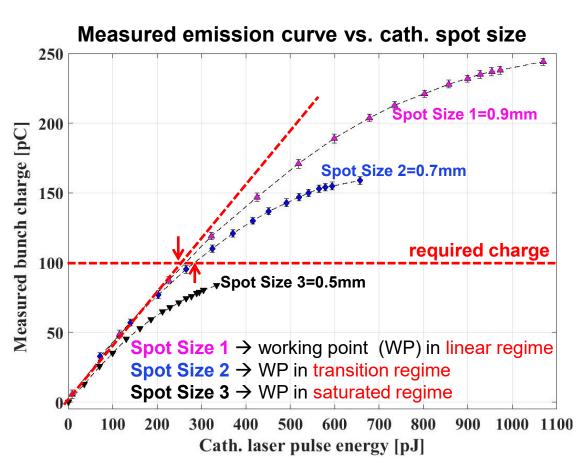
IPAC'19 WEPTS013

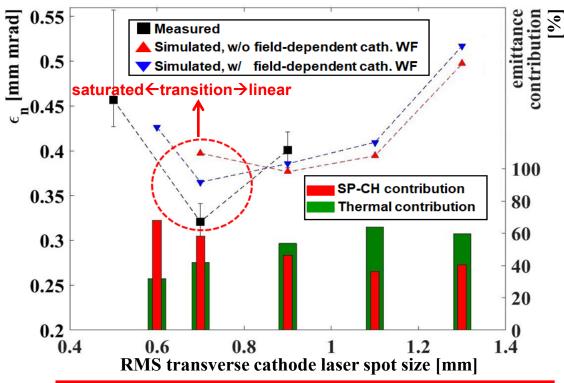


Tracking accelerated bunches (~19 MeV/c) downstream till ~5.3m



Measurement vs. Simulation: optimized emittance vs. cathode laser spot size





- Interplay between space charge emittance and intrinsic emittance gives optimum spot size for best emittance in transition regime
- Improved simulation suggests optimum spot size same as measured

Summary

- Working at transition regime of emission delivers best experimentally optimized emittance
- Photoemission details influence downstream beam qualities
- Emission modeling helps better understand beam dynamics
- Cathode physics important for better emission modeling
- More detailed modeling approach needed for strong space charge fields at cathode

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and many others...

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