

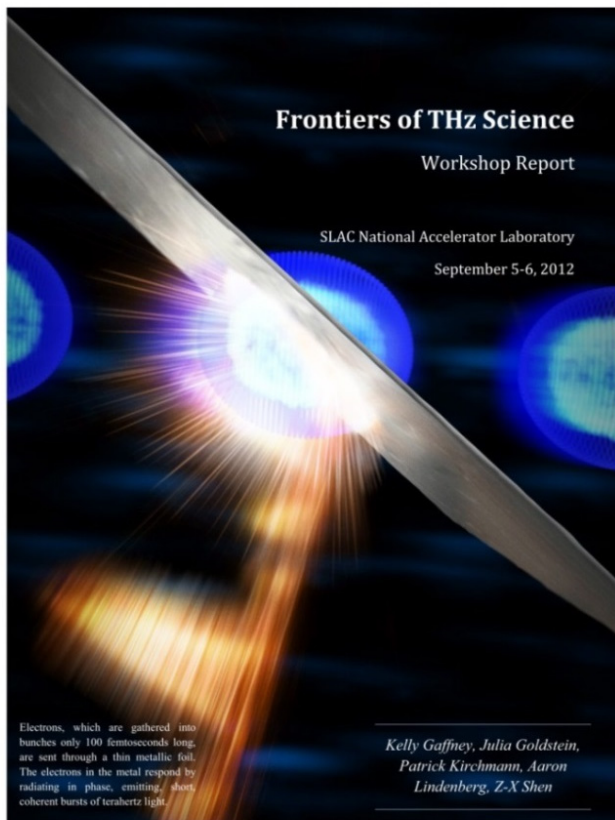


High-power, narrow-bandwidth THz generation using laser-electron interaction in a compact accelerator

Z. Huang, K. Kan, G. Marcus, Z. Zhang (SLAC, Osaka U., Tsinghua U.)

Introduction

- THz science is a rich field with the potential to advance research in many scientific areas.
- Very strong interests exist in combining THz radiation with X-ray FELs for pump-probe experiments



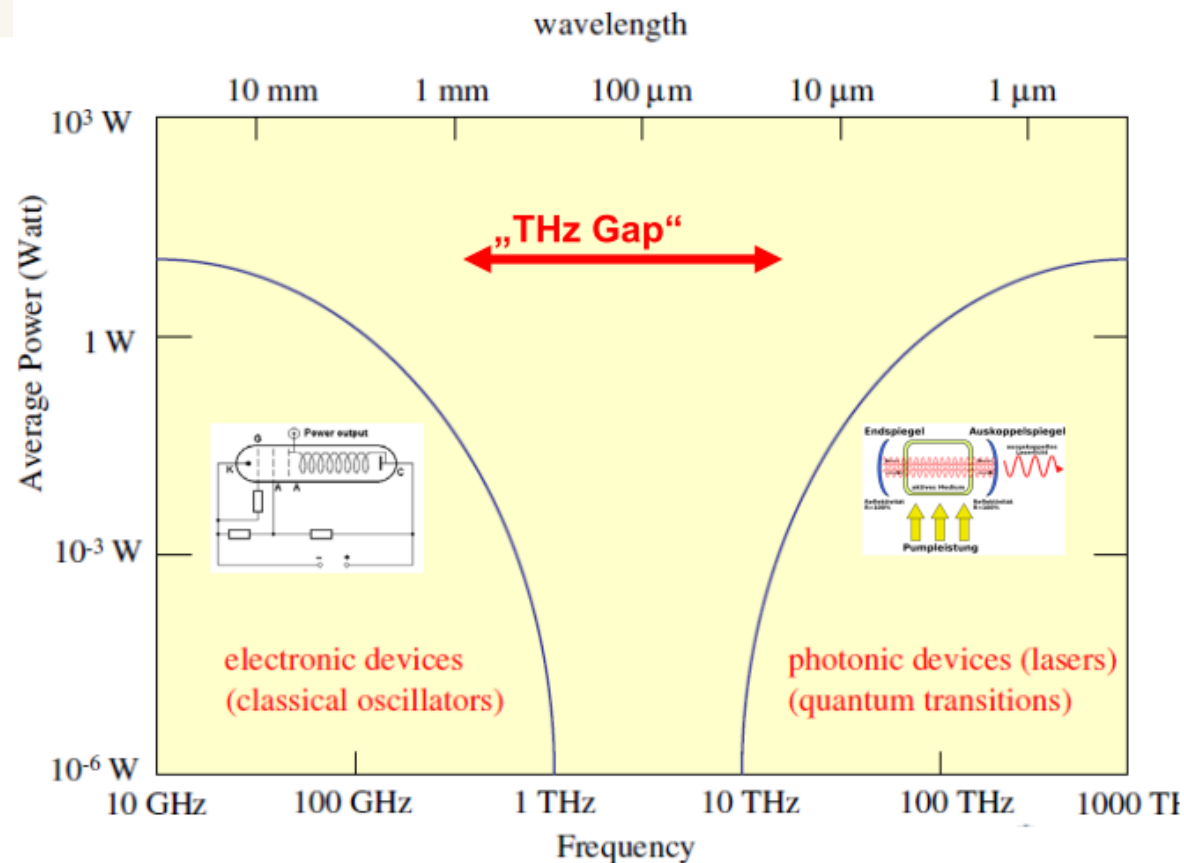
European XFEL Workshop

Terahertz science at European XFEL



01–02 June 2017 / European XFEL, Schenefeld, Germany

THz gap (1 to 20 THz)

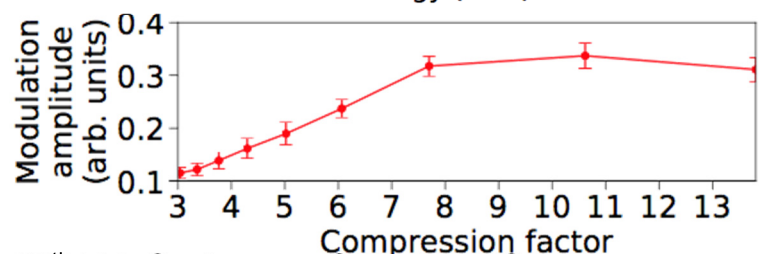
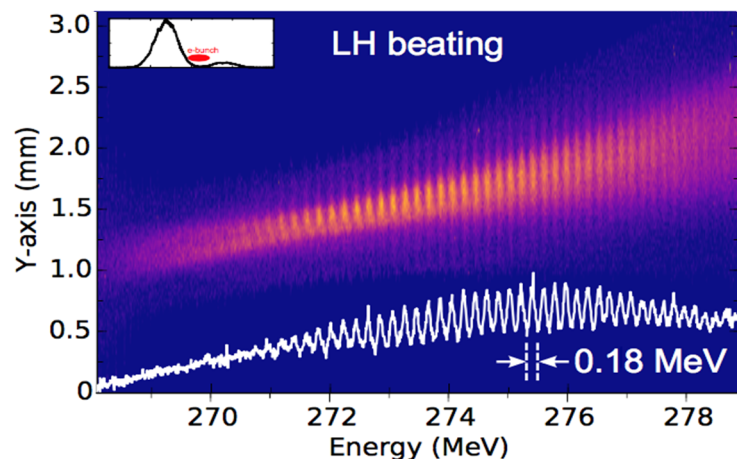
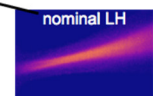
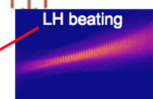
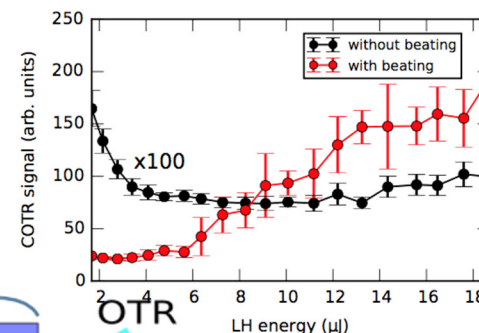
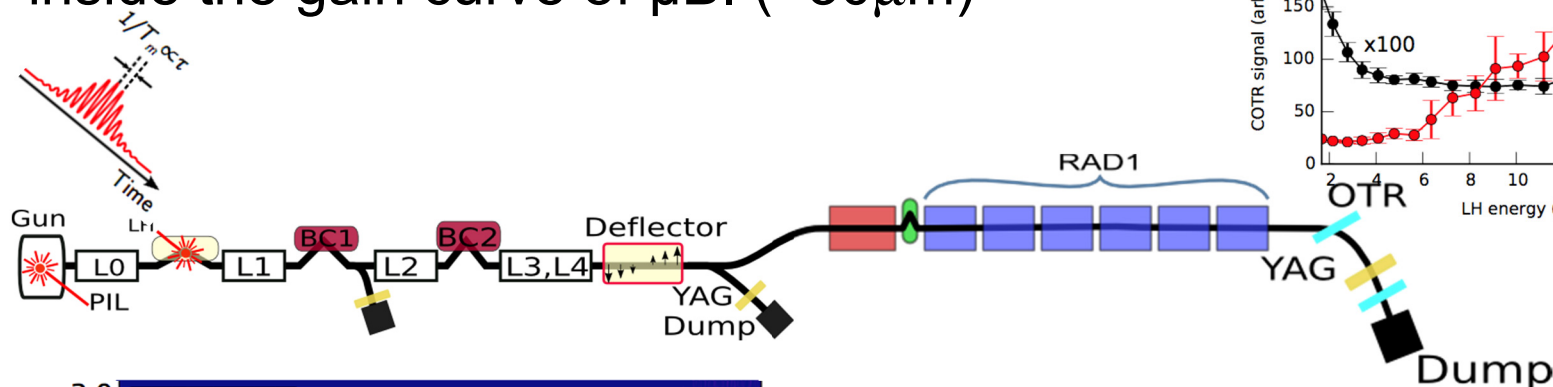


- Laser-based sources have made significant progress, but very challenging to reach above a few THz.
- Accelerator based sources are well suited for high-field, high-frequency, high-rep. rate THz applications

LH beating for seeded μ BI



LH laser adjusted for a beating frequency inside the gain curve of μ BI ($\sim 30\mu\text{m}$)



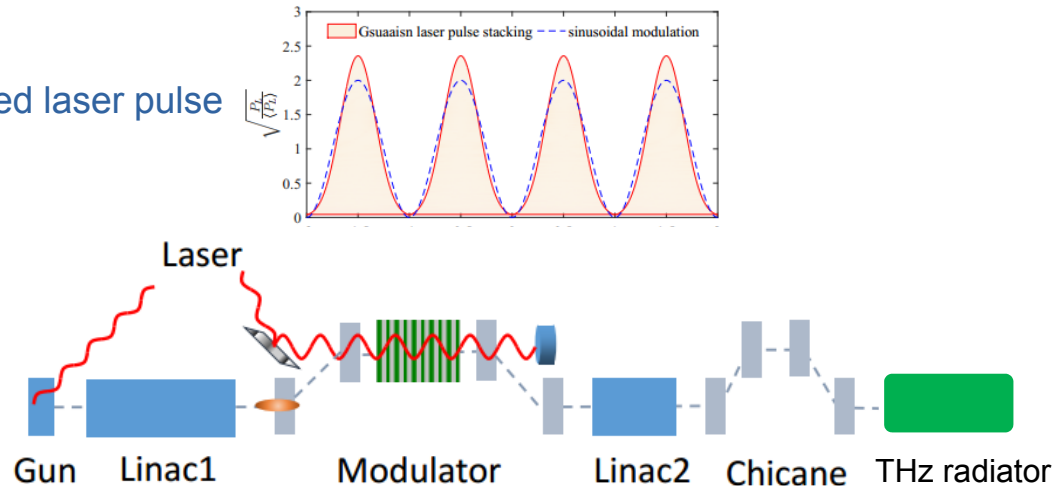
- LH induced **modulation** appears in the e-beam **energy spectrometer** after BC1.
- Spectrum **modulation increases** as the **compression** factor is increasing.
- **COTR** signal measured at the end of the undulator **confirms** that LH induced structures are **transported** to the whole machine.

Our method (Z. Zhang et al., Phys. Rev. AB 20, 050701, 2017)

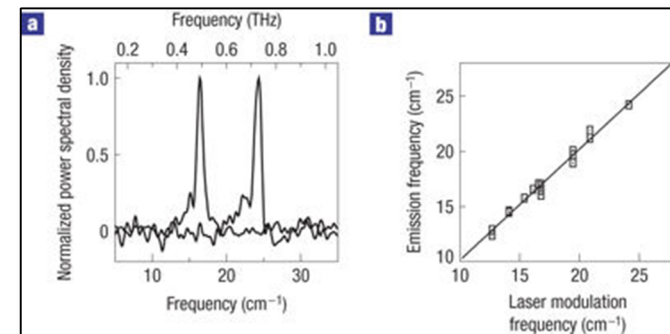
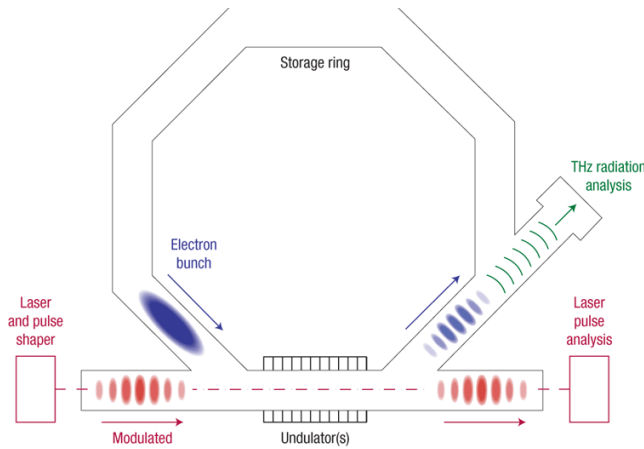
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- We propose a method based on the slice energy spread modulation to generate density bunching in a relativistic electron beam (a la **laser heater setup**)

Amplitude-Modulated laser pulse



- Similar method has been used in storage ring THz generation

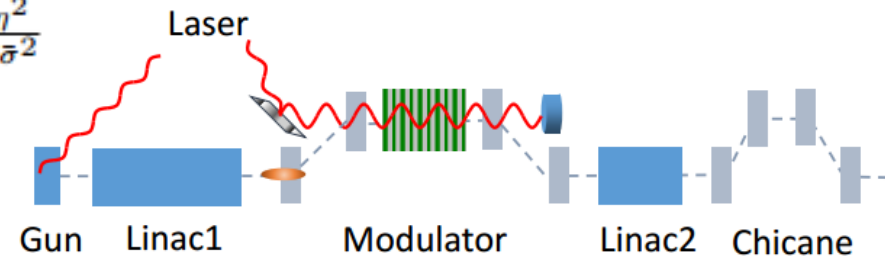


Theory

- We analyze how energy spread modulation is converted to density modulation
- For a beam with gaussian energy spread σ , we find the bunching factor

$$b_n = \frac{(-1)^n}{\sqrt{2\pi\bar{\sigma}}} \int d\eta J_n(k_n R_{56} A \eta) e^{-ik_n R_{56} \eta - \frac{\eta^2}{2\bar{\sigma}^2}}$$

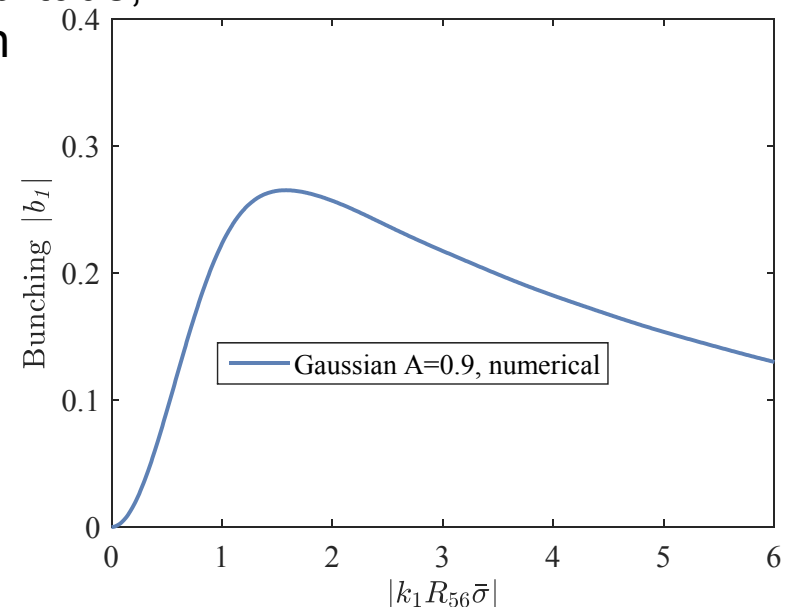
$$k_n = nk_0 / (1 + hR_{56})$$



n harmonic #, k_0 initial mod. wavenumber, A is amplitude,
 h electron energy chirp, R_{56} is the chicane strength

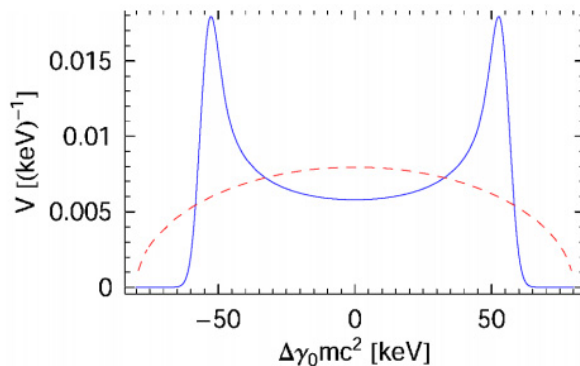
- For the first harmonic $n = 1$, the optimal chicane setting is

$$|k_1 R_{56} \bar{\sigma}| \approx 1.75$$



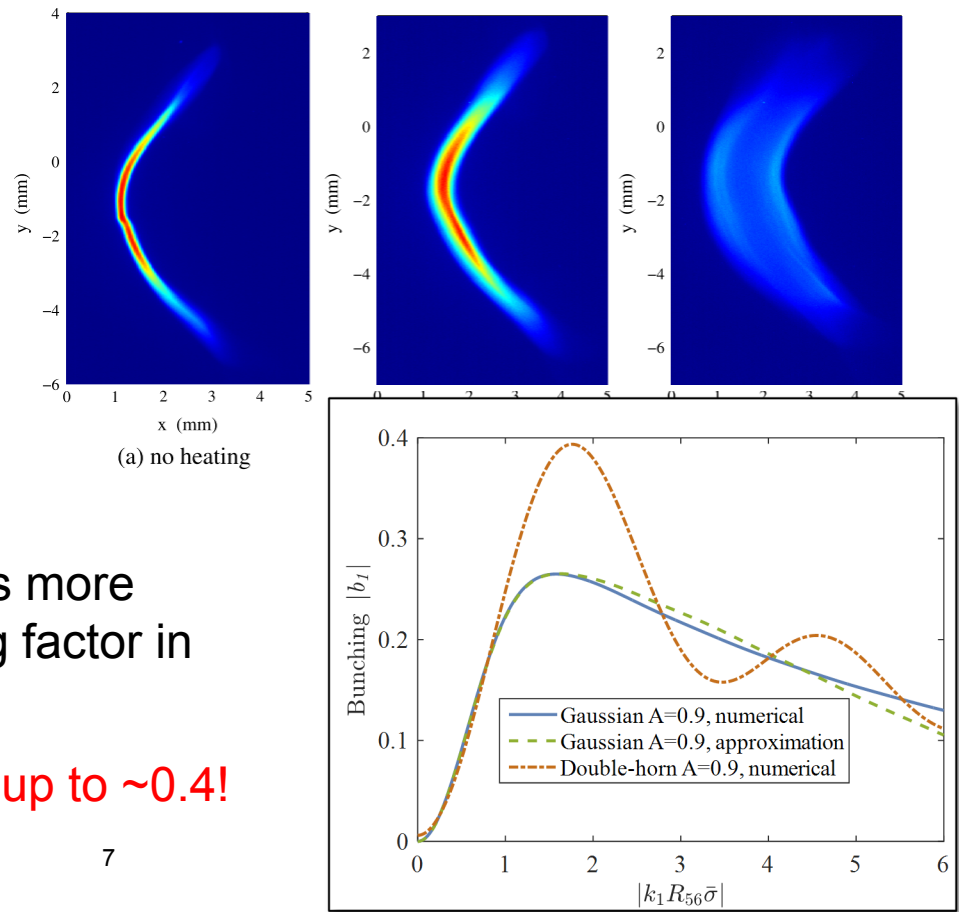
Theory

- The analysis above assume the beam has a Gaussian slice energy distribution, which is not true for the laser electron interaction.
- When the laser waist size in the undulator is much larger than the beam size, the resulting energy profile is a double-horn distribution



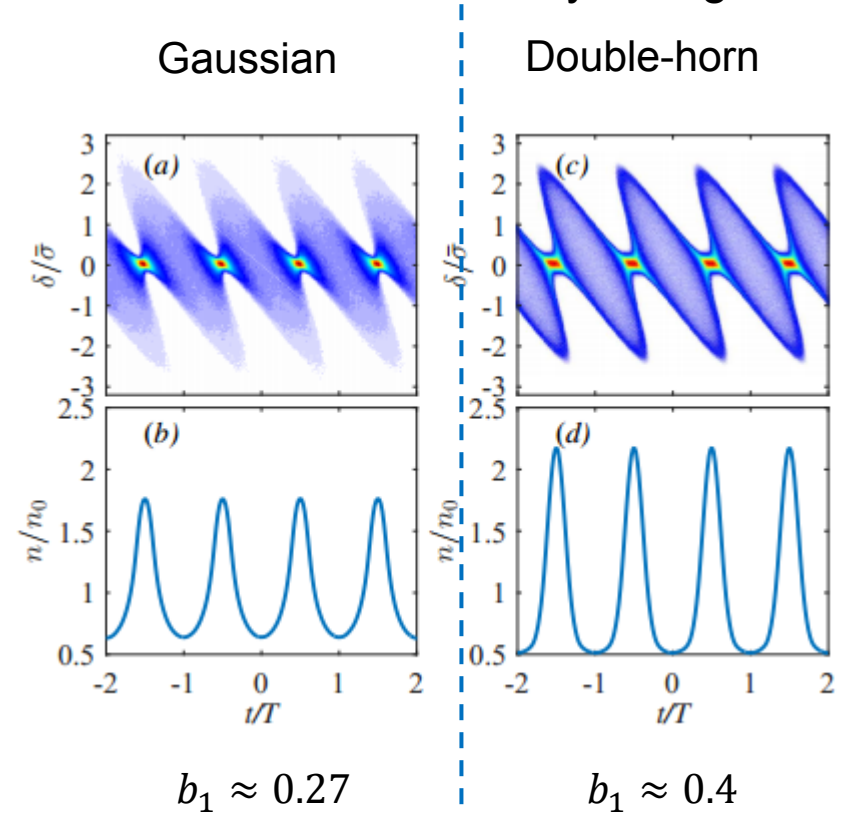
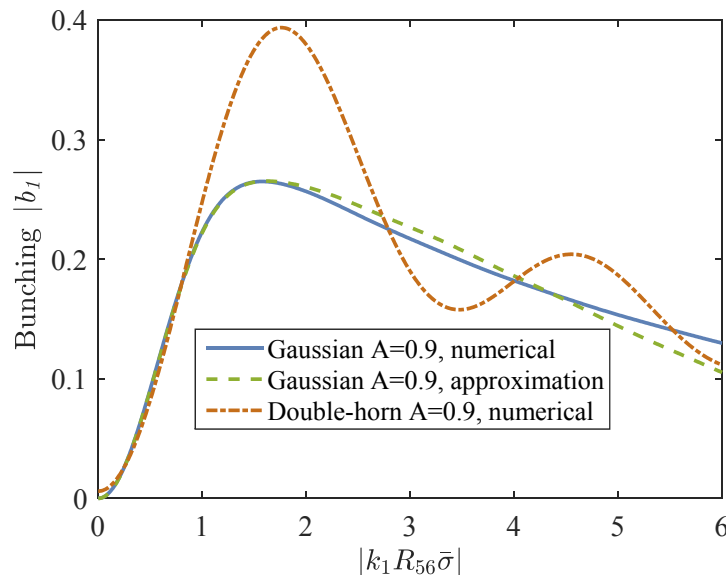
Z. Huang et al., PRST-AB 7, 074401 (2004)
PRST-AB 13, 020703 (2010)

- Double-horn energy distribution is more effective to increase the bunching factor in our study.
- The maximum bunching factor is up to ~ 0.4 !



Theory

- Phase spaces of Gaussian and double-horn distributions when yielding maximum bunching factor



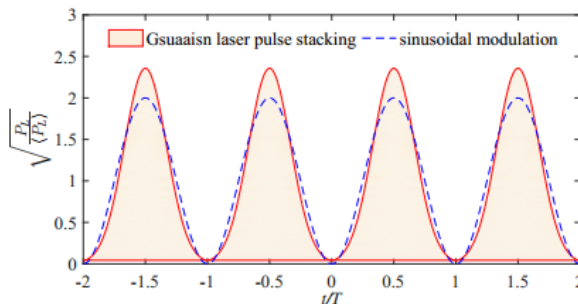
- Significant second harmonic bunching (~ 0.25) can be used to reach higher THz frequencies

Simulation setup

- Parameters similar to LCLS or FERMI injector
- The laser modulation can be generated by the chirped pulse beating or pulse stacking

Parameter	Value	Units
Electron beam		
Charge	500	pC
Beam energy	135	MeV
Current Profile	flat-top	/
Bunch length	~ 10	ps
Intrinsic slice energy spread	10^{-4}	/
Norm. emittance	1	mm-mrad
rms beam size	200	μm
Modulator		
Laser wavelength	800	nm
Undulator period	5	cm
Period number	10	/
Laser waist size	1.5	mm
Laser stacking separation	0.5 (0.25*)	ps
rms laser pulse length	60 (30)	fs
Laser power	1 (0.26)	GW

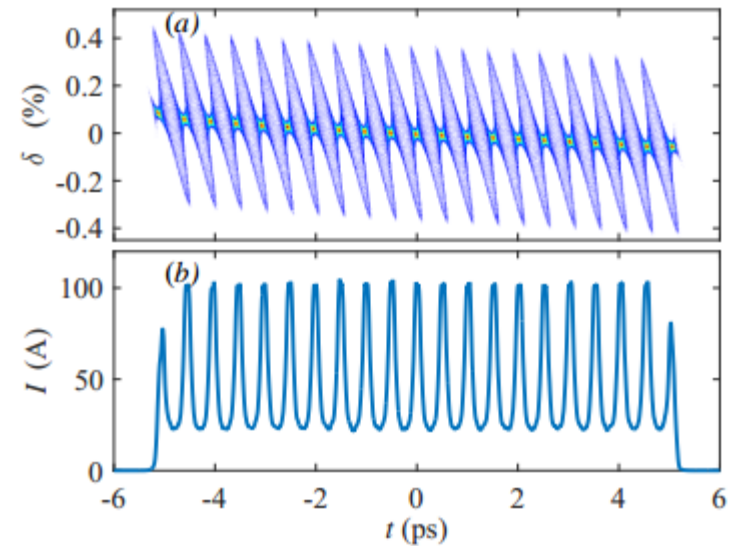
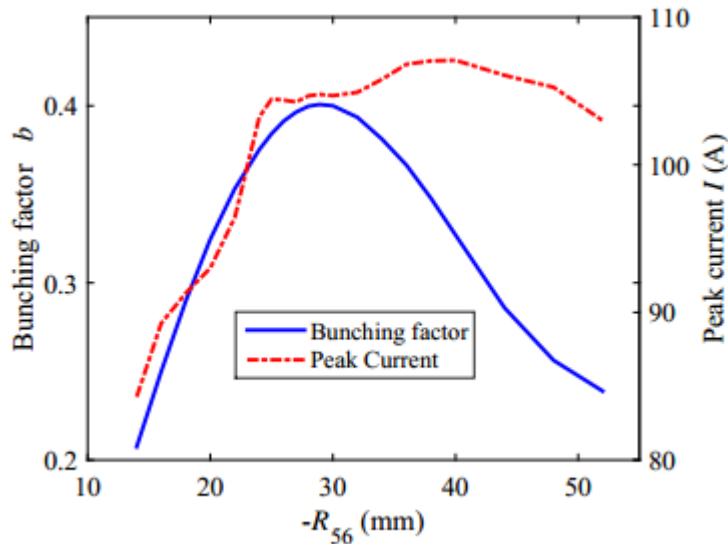
* The numbers in brackets are the parameters for 4 THz case.



period $T = 0.5\text{ps}$ (0.25ps)
frequency $f_0 = 2\text{ THz}$ (4 THz)
 $\sigma_t = 60\text{fs}$ (30 fs)

Elegant simulation results

- 2THz initial modulation, scan R_{56} , parameters: $P = 1\text{GW}$, $\bar{\sigma} = 190\text{keV}$ (1.4×10^{-3})



- The optimal condition $|k_1 \bar{\sigma} R_{56}| \approx 1.75$ predicts the optimal chicane is -29.4mm, consisting with the simulations (-29mm).
- The peak current stays constant with larger R_{56} .
- Similar results for 4 THz initial modulation. Electron energy chirp can be used to tune density modulation frequency from 2-6 THz.
- The longitudinal phase space and current profile when the $R_{56} = -29\text{mm}$ (optimal bunching)

Stand-alone compact accelerator-based THz source

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- Compact accelerator of ~50 MeV interacts with a 800-nm laser in the undulator
- Laser-electron interaction through 3rd harmonic (for a planar undulator with fundamental resonant wavelength 2.4 μm .)

Parameter	Value
E-beam charge	1 nC
Beam energy	50 MeV
RMS beam size	0.2 mm
Bunch length (flat-top)	10 ps
Modulator	
Undulator period	2.5 cm
Peak field/ K value	0.56 T / 1.29
Undulator length/period	0.5 m / 20
Laser wavelength	800 nm
Laser RMS spot size	0.5 mm
Laser stacking separation	0.5 ps (2.7 THz)
Laser peak power	100 MW

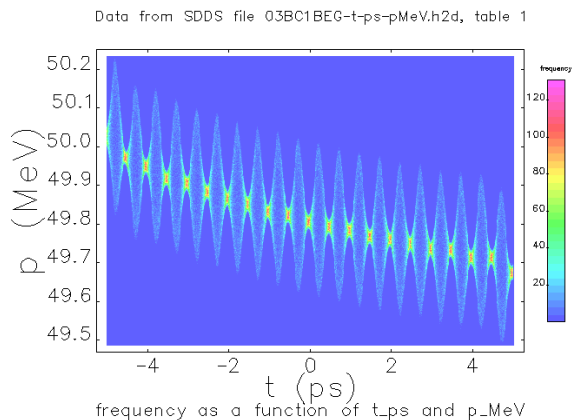
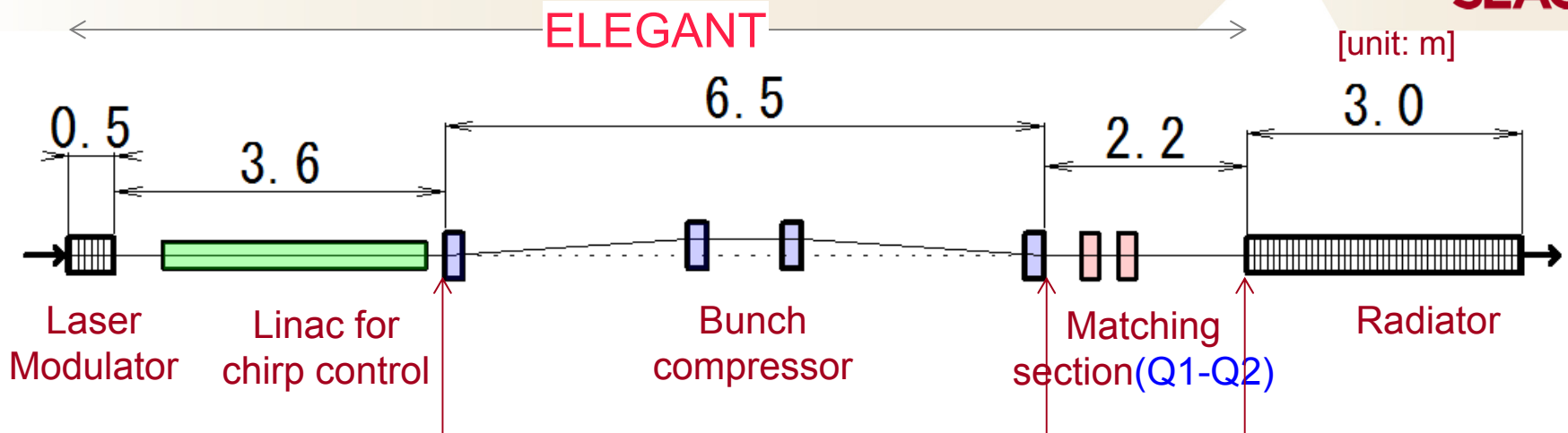
Relative energy modulation

$$\delta = \sqrt{\frac{P_L}{P_0}} \frac{KL_u}{\gamma^2 \sigma_r} [JJ]_3$$

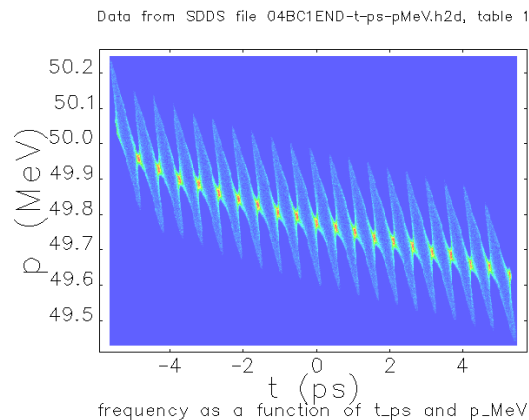
OPCPA laser at high-rep. rate
(100 kHz, 100 W average power)

50 MeV simulation results

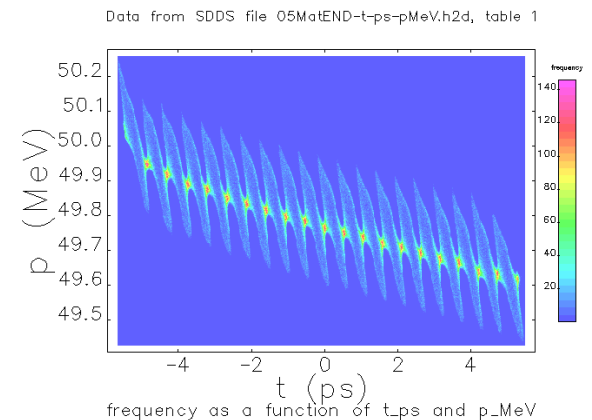
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Before bunch compressor
(2 THz modulation
rms Espread=0.15%)



After bunch
compressor
(bunching=0.37)



After matching
section
(bunching=0.34)

THz radiation

- Transition radiation foil is the simplest radiator. CTR energy ~ 10 μJ (5% bandwidth) with 1 nC charge.
- For a helical wiggler ($\lambda_w = 15$ cm, $K_w = 4.3$), resonant wavelength at 50 MeV is $\lambda_r = 150 \mu\text{m}$ (2 THz).
 - Radiation pulse energy for a thin beam (large diffraction regime)

$$W_0 = W_b \left[\frac{\pi^2 a_{\text{in}}^2}{2} \right] \left[\frac{I}{\gamma I_A} \right] \left[\frac{K_w^2}{1 + K_w^2} \right] N_w$$

Saldin, Schneidmiller, Yurkov, NIMA539, 499 (2005)

W_b is the beam power (50 mJ at 1 nC), $a_{\text{in}} = 2b = 0.68$, $I = 100$ A,

$N_w = 20$ (3 m wiggler) $\rightarrow$ THz pulse energy $W_0 = 140 \mu\text{J}$

- Undulator in a waveguide or a dielectric tube may be more efficient radiators

Summary

- Based on the slice energy spread modulation method, the bunching factor can be kept at ~ 0.4 for a wide frequency range (1-10 THz) and can be tuned by laser amplitude modulation frequency or electron energy chirp.
- THz pulse energy is estimated to be 10-100 μJ for a compact accelerator (40~50 MeV). Further optimization can yield mJ pulse energy.
- THz pulse shape can be controlled by both laser and e-beam techniques (narrowband, chirped, a few cycle pulses, all possible). Rep. rate in the kHz-MHz range is achievable.
- THz modulation experiments are planned at FERMI FEL as well as on a compact accelerator in Tsinghua University.

Acknowledgement

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