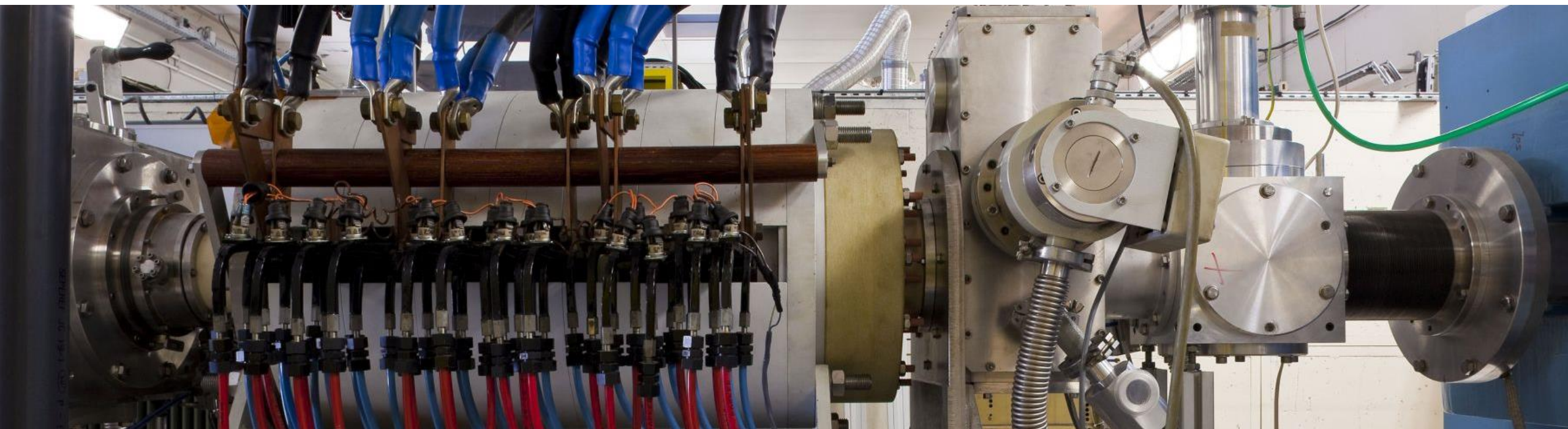


Cyclotron Ion Sources

The capture of injected 1+ ions into an ECRIS CB plasma

J. Angot¹, O. Tarvainen², P. Chauveau³, A. Galatà⁴, L. Maunoury^{3*}, T. Thuillier¹



For more details on this R&D :

OPEN ACCESS
IOP Publishing


Plasma Sources Science and Technology







Plasma Sources Sci. Technol. **31** (2022) 125003 (7pp)


<https://doi.org/10.1088/1361-6595/aca713>

The capture of injected 1+ ions in charge breeder electron cyclotron resonance ion source plasma

RESEARCH ARTICLE | MAY 06 2024

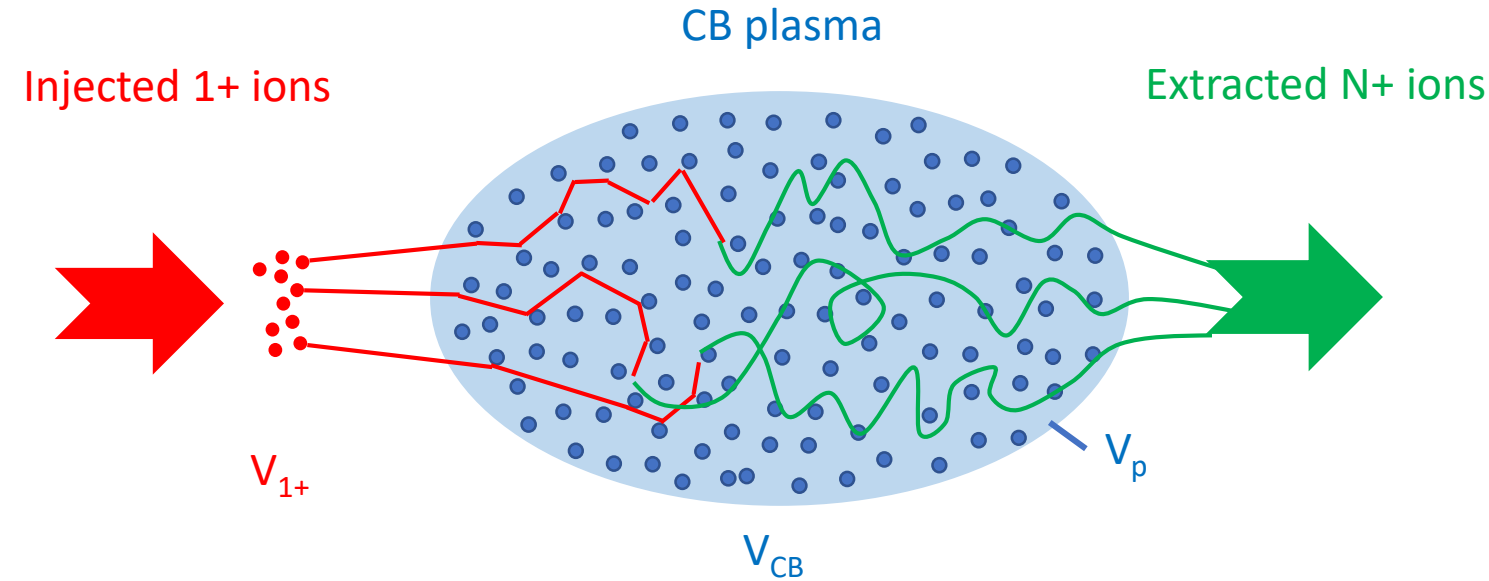
Experimental investigation of electrostatic capture of 1+ ions in charge breeder electron cyclotron resonance ion source plasma 

O. Tarvainen   ; J. Angot   ; P. Chauveau; A. Galatà  ; T. Thuillier 

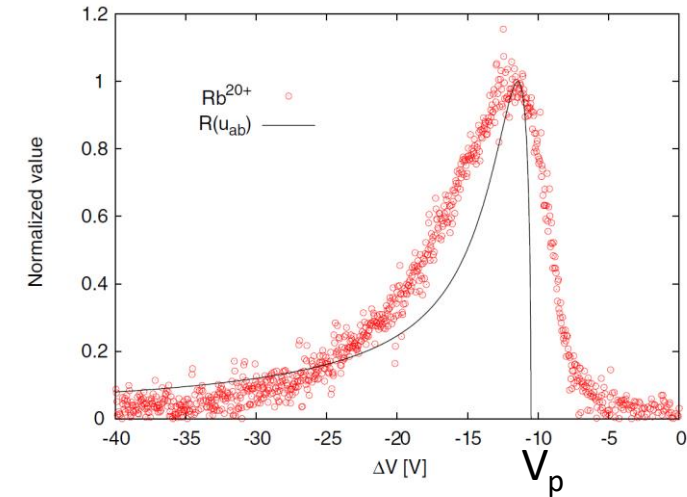


Phys. Plasmas **31**, 053103 (2024)

<https://doi.org/10.1063/5.0202875>



Rb²⁰⁺ ΔV curve, O₂ as support gas



Considering a Maxwellian velocity distribution for the plasma ions :

$$S_{ab} = \frac{\langle \Delta v_{a,||} \rangle}{\Delta t} \quad S_{ab} = -\frac{n_b}{4\pi\epsilon_0^2} \left[\frac{q_a q_b e^2}{m_a \langle v_b \rangle} \right]^2 \left(1 + \frac{m_a}{m_b} \right) R(u_{ab}) \ln \Lambda$$

$$R(u_{ab}) = \frac{2}{\pi^{1/2}} \frac{1}{u_{ab}^2} \int_0^{u_{ab}} x^2 e^{-x^2} dx \quad u_{ab} = \frac{v_{a,||}}{\langle v_b \rangle}$$

$$R(u_{ab}) \text{ maximum when } u_{ab}=1 \text{ so when } v_{a,||} = \langle v_b \rangle$$

a injected species, b plasma species

n density, q charge, m mass, v velocity, $\ln \Lambda$ Coulomb logarithm

$$E_{\text{opt},1+} = \frac{4m_{1+}}{\pi m_i} kT_i$$

$$\Delta V = V_{CB} - V_{1+}$$

$$E_{1+} = e(|\Delta V| - V_p)$$

$$|\Delta V|_{\text{opt}} = \frac{4m_{1+}}{\pi m_i} \frac{kT_i}{e} + V_p$$

$$|\Delta V|_{opt} = \frac{4m_{1+}}{\pi m_i} \frac{kT_i}{e} + V_p$$

From experimental results :

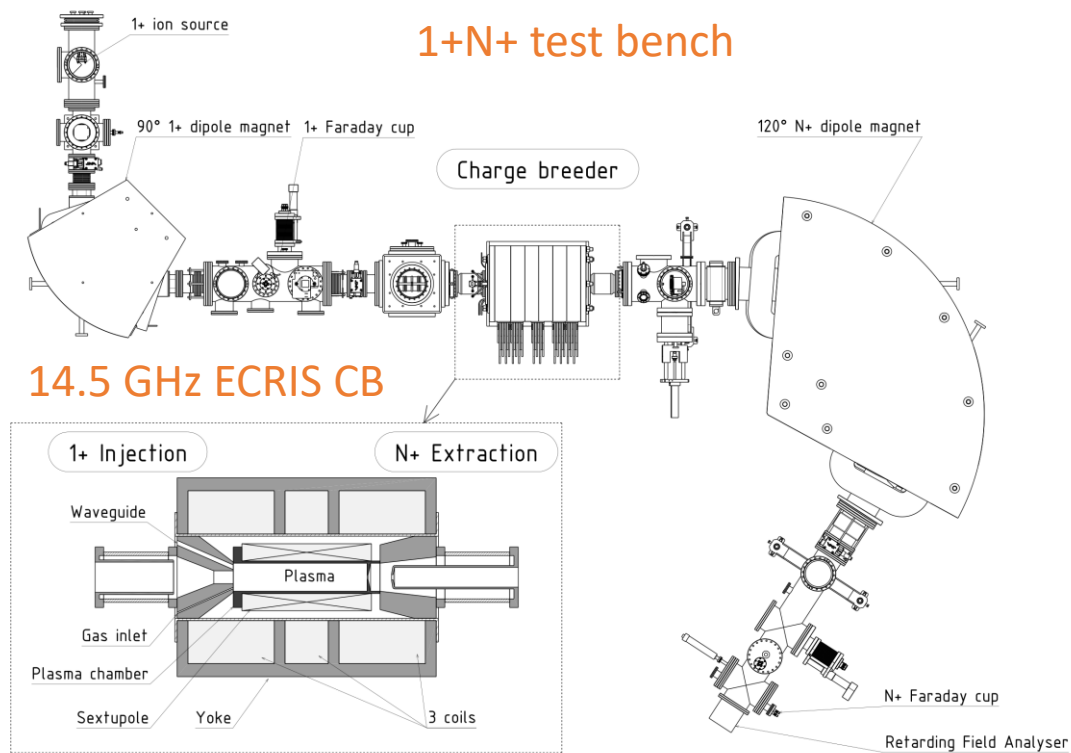
- $m_{1+}/m_i > 5$ for good ECR CB efficiency
 - Ion temperature in ECRIS plasma 5 - 28eV (optical spectroscopy measurements)
- $|\Delta V|_{opt} > 50V$ for Na^+ injected into an He plasma, even considering $V_p = 0V$
- 5 – 15 V values are typically measured

We propose here a revised concept where the slowing down is mainly electrostatic due to V_p

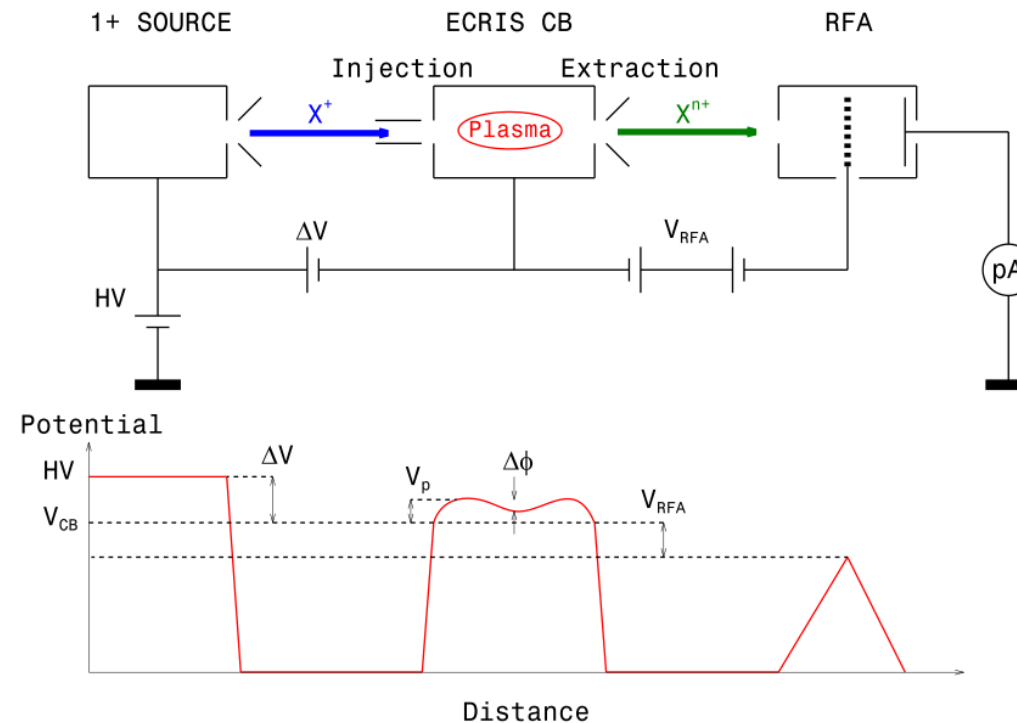
$$E_{opt,1+} = eV_p$$

We did extensive experiments to verify this assumption.

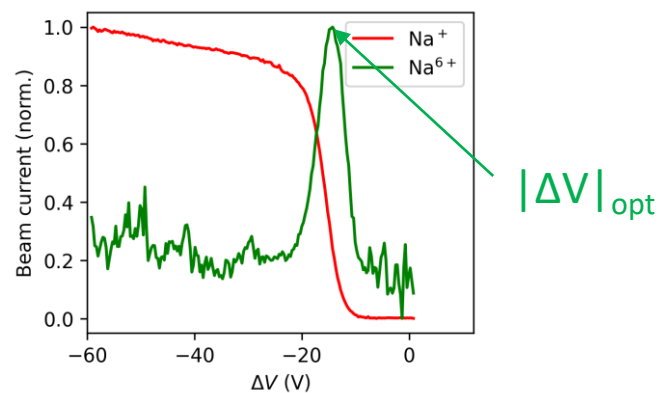
Experimental setup



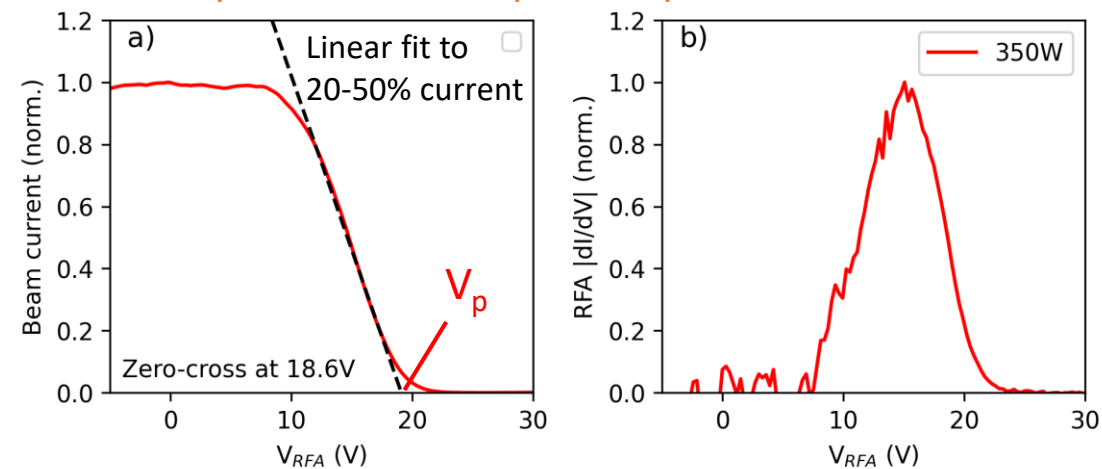
Electrical configuration



ΔV spectrum

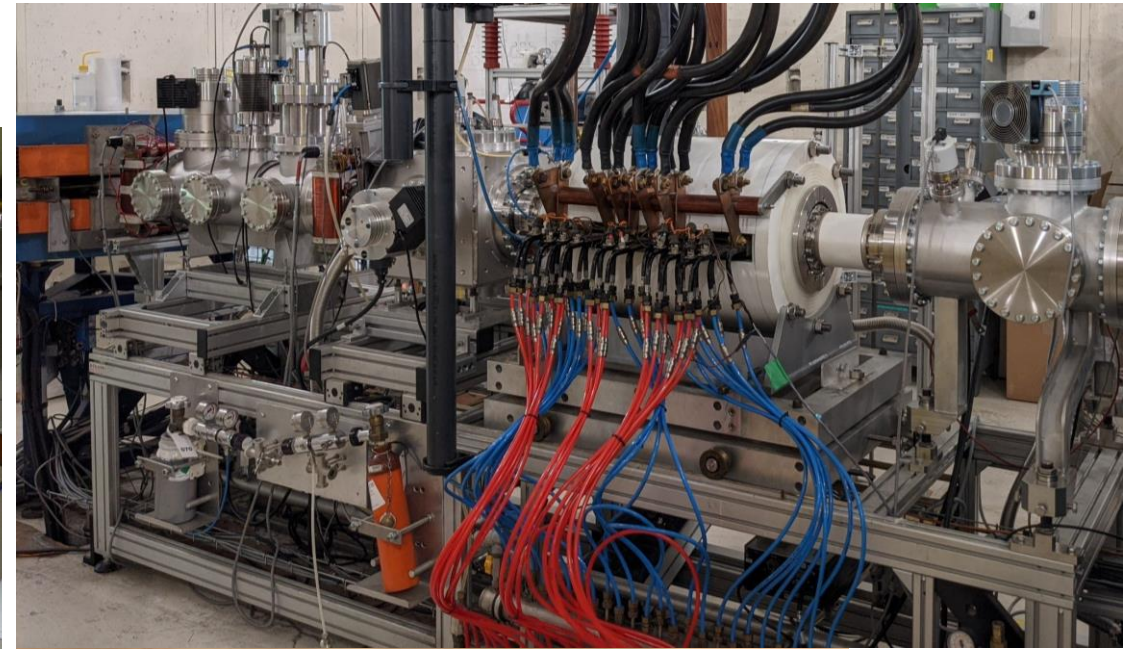
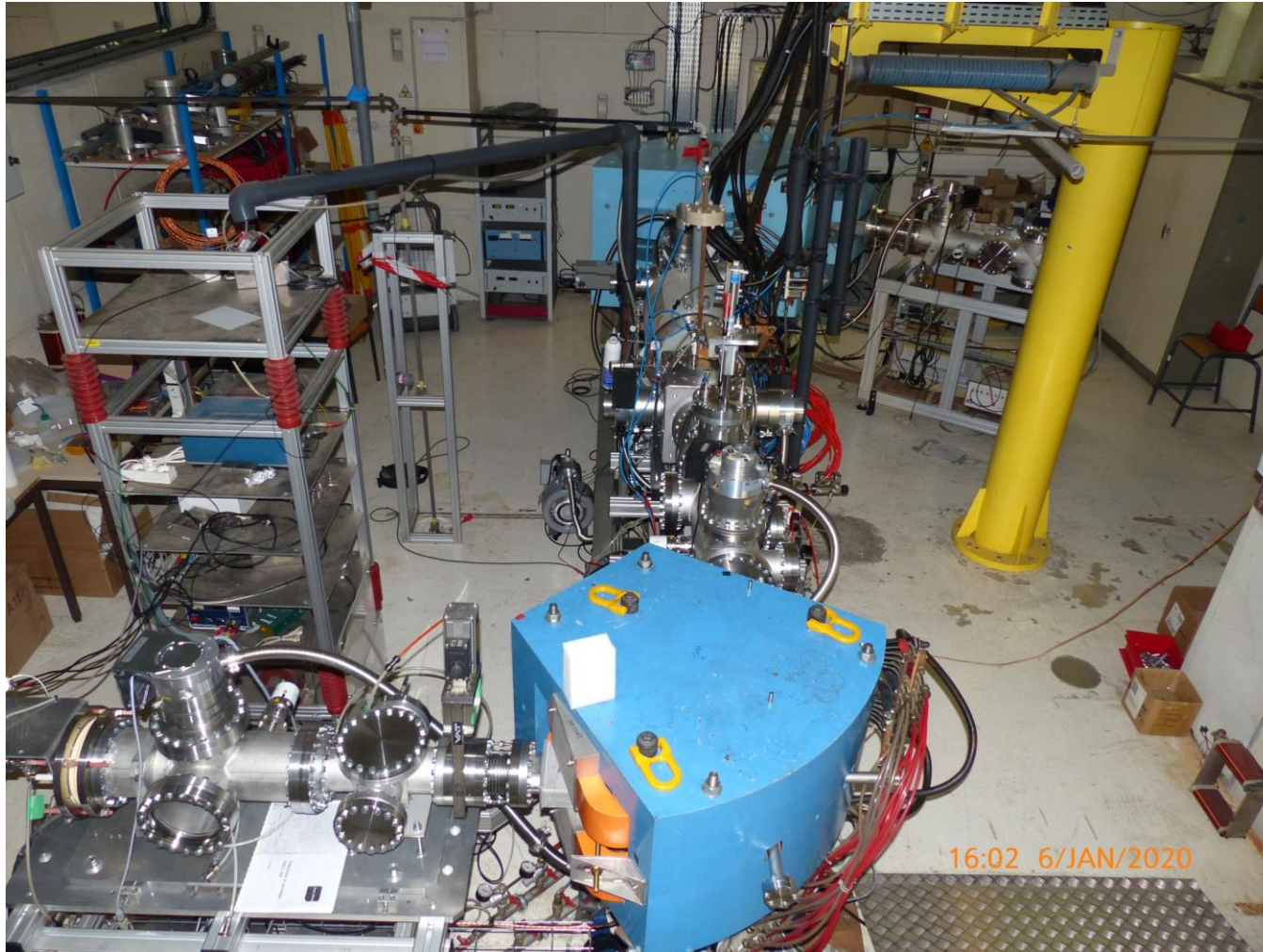


RFA spectrum , with plasma potential assessment



Experimental setup

1+N+ test bench

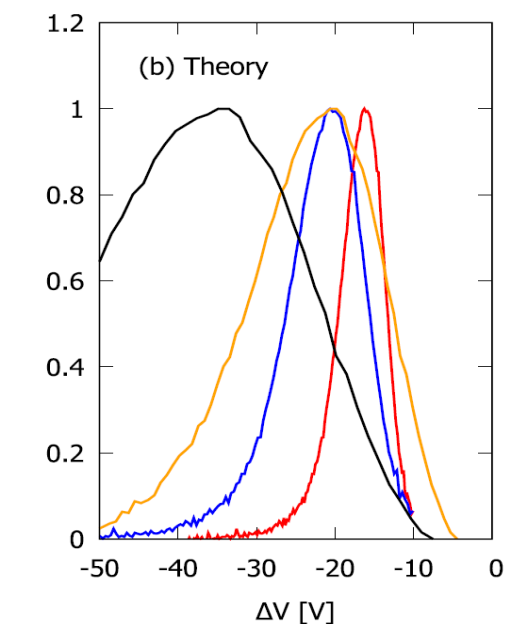
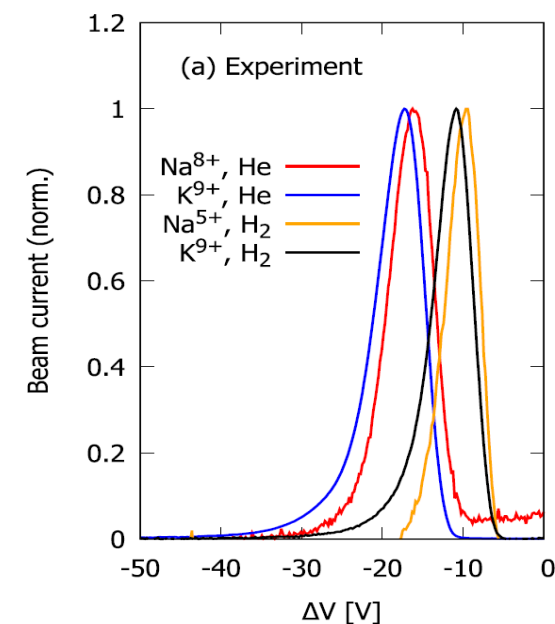
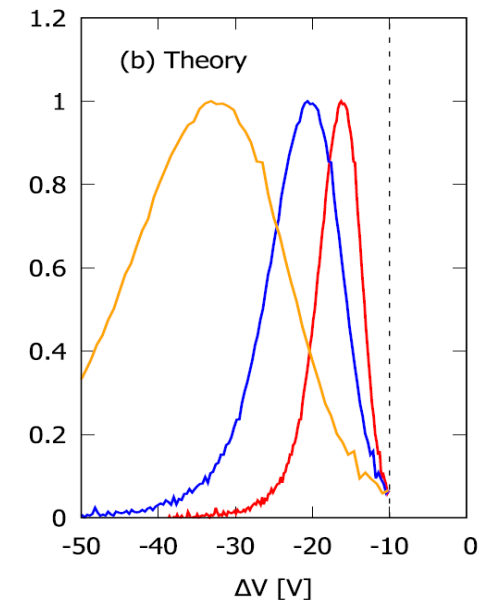
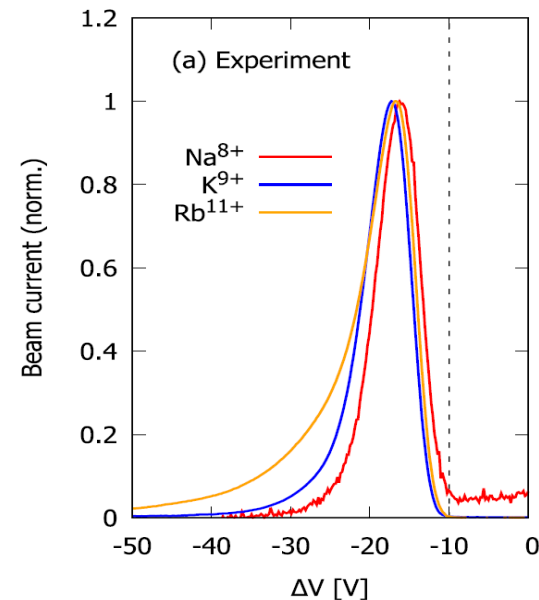


Check the mass dependency of $|\Delta V|_{\text{opt}}$:

$$|\Delta V|_{\text{opt}} = \frac{4m_1}{\pi m_i} \frac{kT_i}{e} + V_p$$

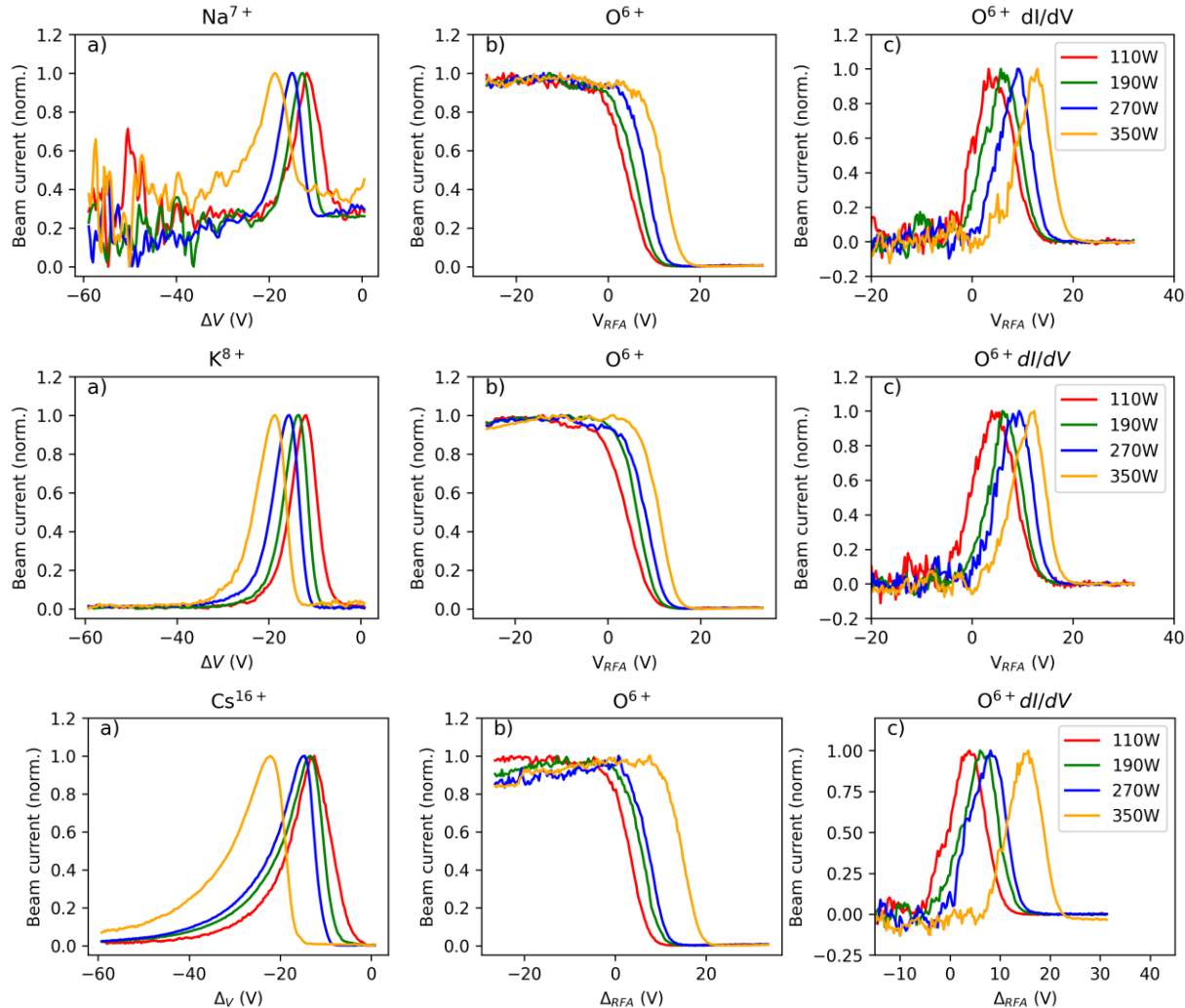
Mass ratio

He plasma support gas



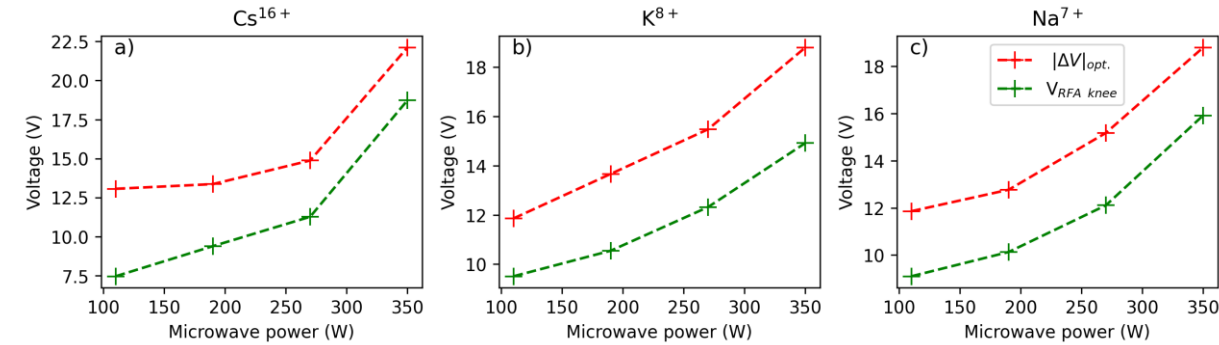
$$|\Delta V|_{opt} = \frac{4m_1 + kT_i}{\pi m_i} \frac{1}{e} + V_p$$

Plasma potential variation with microwave power



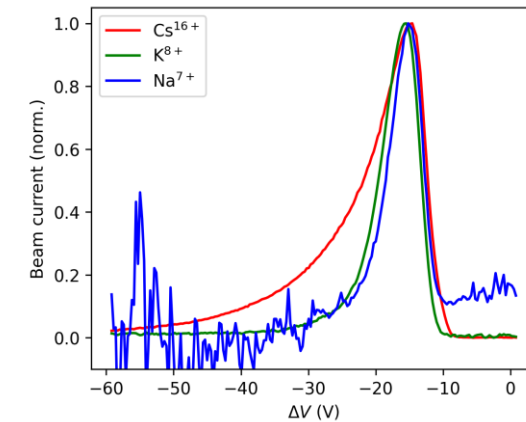
Na+ and K+ injected exactly in the same plasma conditions

Plasma potential and $|\Delta V|_{opt}$ vs microwave power



➤ $|\Delta V|_{opt}$ follows the plasma potential with an offset

DeltaV curves comparison at 270W microwave power

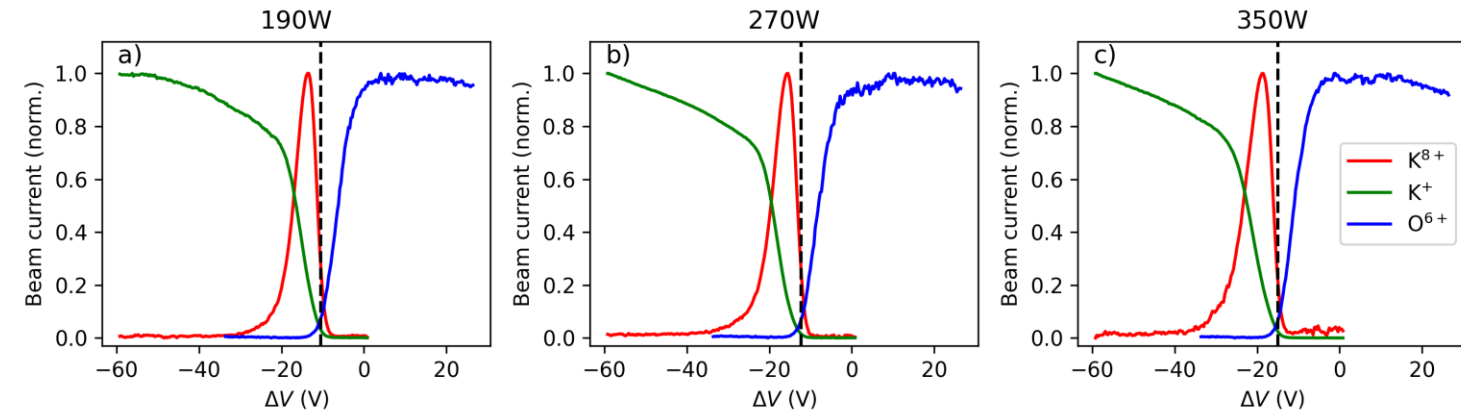


- Same $|\Delta V|_{opt}$ whatever the injected mass
- Larger FWHM for Cs : lower for high mass ?

Results

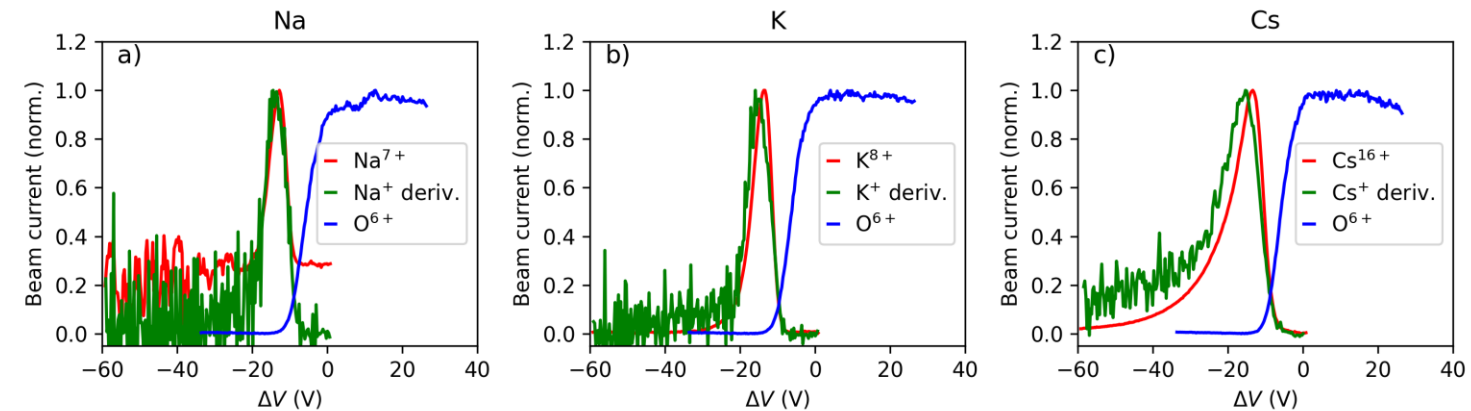
ΔV spectrum of K at different microwave power
With RFA spectrum (inverted voltage) to estimate the plasma potential value (dashed lines) :

- When the energy of injected ions becomes greater than $e.V_p$, they can propagate through the CB plasma.



ΔV spectrum of Na K Cs at 190W microwave power, derivative of 1+ (fly through) charge state :

- Capture optimum corresponds to the most efficient stop of the 1+ ions
- The 1+ ions stopped are effectively converted to high charge state ions in the energy range
- Small shift of optimum visible : in flight ionisation to 2+

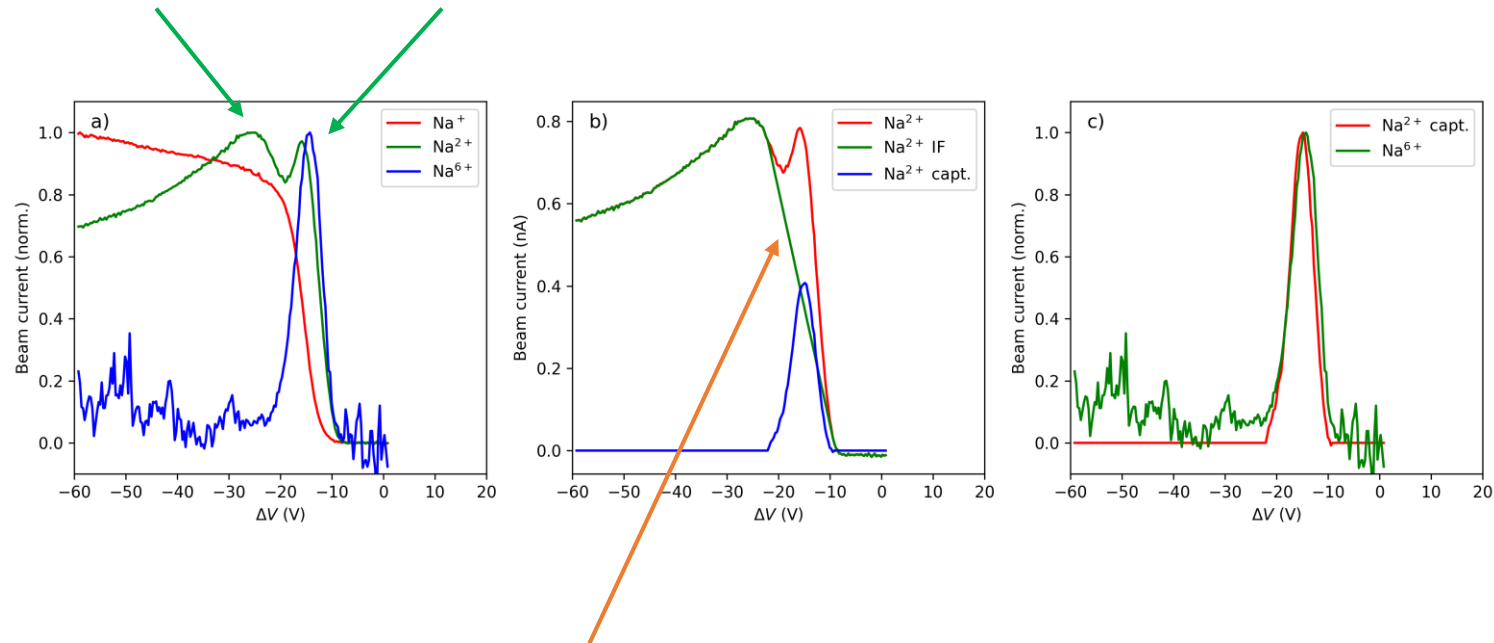


In flight $1+ \rightarrow 2+$ ionisation

Optimum fly through $1+ \rightarrow 2+$ ionisation

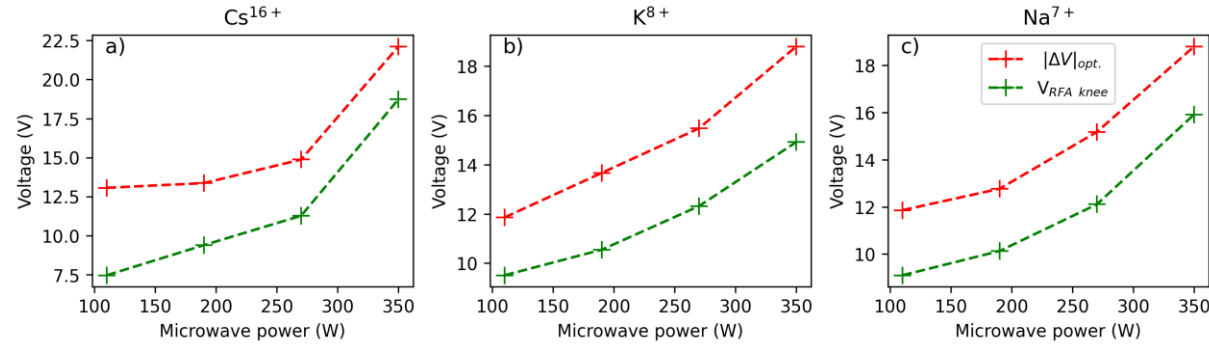
Optimum captured $1+ \rightarrow 2+$ ionisation

Na^+



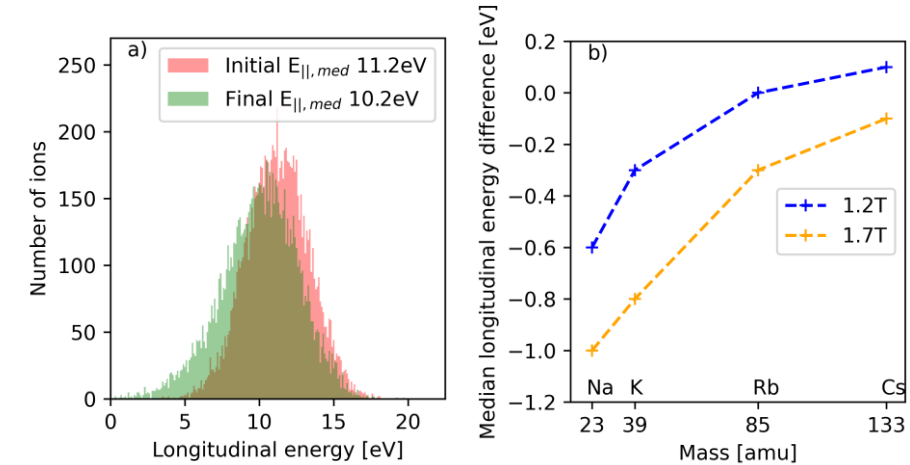
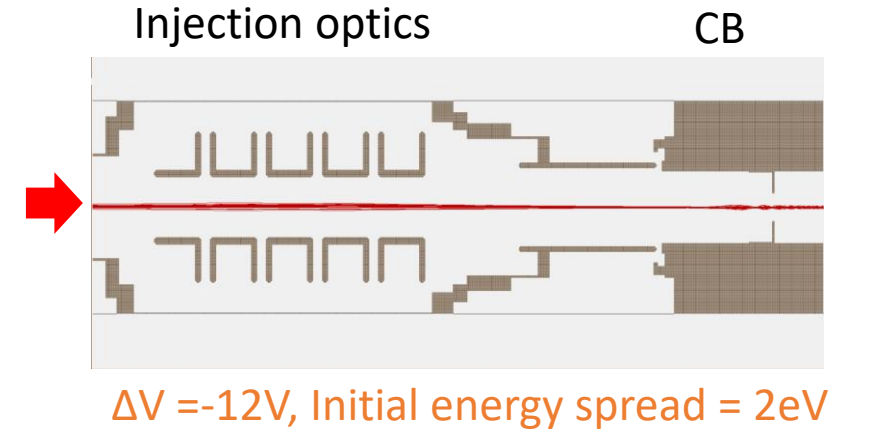
Higher fly through $1+ \rightarrow 2+$ ionisation at higher energy could lead to « shift » the $|\Delta V|_{\text{opt}}$ for high charge state at lower energy

Discussion



Possible explanation of the difference between $|\Delta V|_{\text{opt}}$ and plasma potential value
It is in the range $2.4 \text{ V} < \text{Offset} < 5.6 \text{ V}$

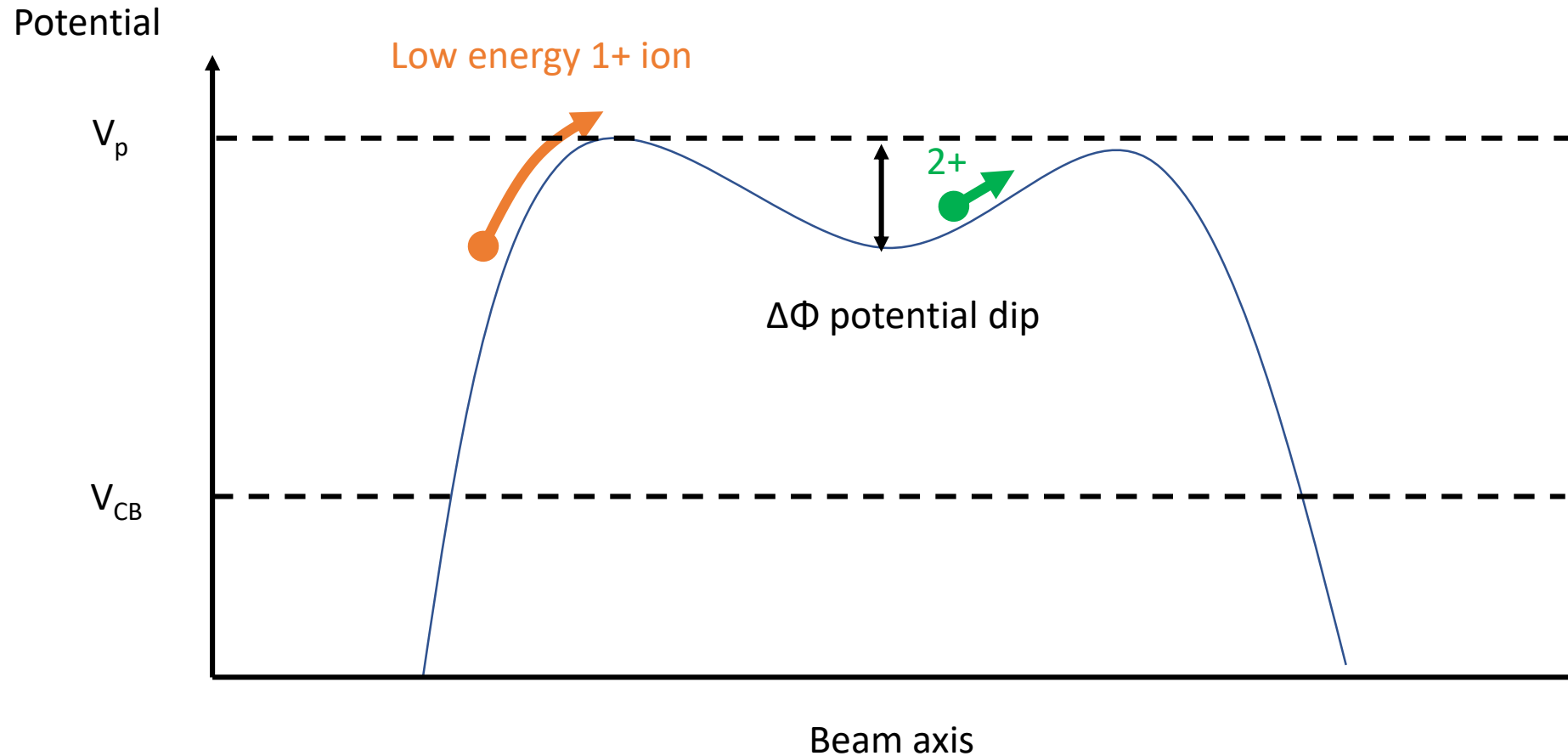
- Injection magnetic field : converts longitudinal energy to transverse
Study of the effect by simulation with Simion
- Plasma potential (peak value) not symmetric at injection vs extraction
- Coulomb interactions before the injected ions reach the plasma potential peak value
- Injection with an excess of energy balancing the effect of inflight ionisation to 2+ and reflection at injection side



Discussion

Electron impact ionisation to 2+

- Enhanced friction (proportional to q^2)
- Electrostatic trapping in plasma potential dip (required energy doubles)



Conclusion

Experiments have demonstrated that the capture of 1+ ions injected into an ECR charge breeder plasma is dominated by the electrostatic forces :

- No direct mass dependency can be observed on the $|\Delta V|_{\text{opt}}$ parameter
- $|\Delta V|_{\text{opt}}$ follows the plasma potential value with a small offset

$$E_{\text{opt},1+} = eV_p$$

In the proposed model, the capture doesn't depend on the plasma ion temperature



Thank you for your attention