

Measurement of phase-space density evolution in MICE Step IV

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on behalf of the MICE collaboration

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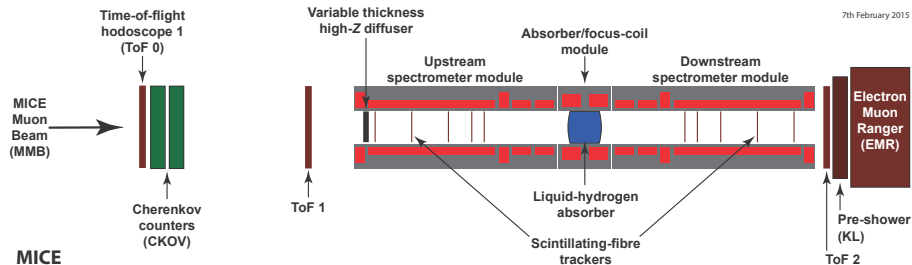
Experimental apparatus at present (Step IV)

All the detectors are installed and **working**

- Three time-of-flight (TOF) detector stations
- Two Cherenkov counters and a downstream calorimetry module
- Two scintillating-fibre trackers

Part of the cooling channel (no RF yet)

- Two Spectrometer Solenoids (SS), each composed of 5 coils
 - ▶ The two downstream match coils are currently **not turned on**
- An Absorber Focus Coil (AFC) module, made of 2 coils + absorber

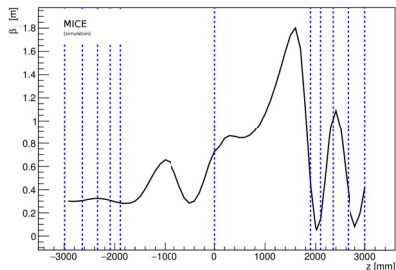
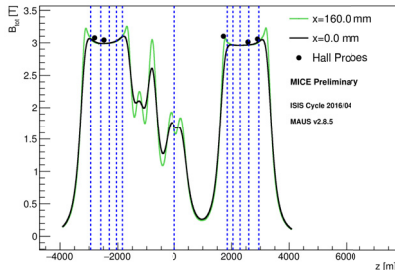


Beam cooling setting optimization

Suitable optics have been found using two approaches to conjointly optimize transmission and cooling performance:

- Linear optics, scan in the parameter space of magnet currents;
- Genetic algorithm, best sets of optics bear the next generation, penalize transmission loss and encourage emittance reduction.

The **bottom lattice** setting is expected to have one of the best cooling performance – transmission trade-off and is **presented in the following**

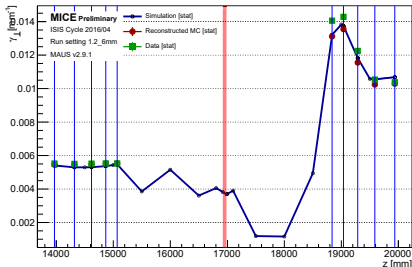
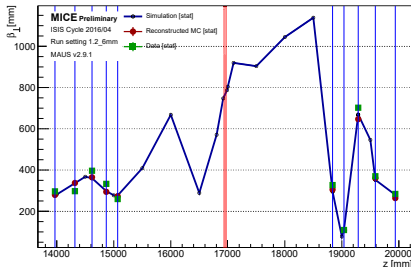
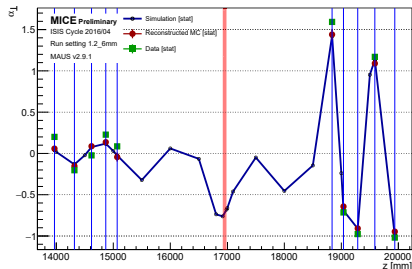


Reproduction of optical functions in the simulation

Reliable reproduction of the data in the simulation for

- ~ 6 mm input beam
- 140 MeV/ c central momentum
- -0.68 central $\alpha_{\perp}$
- 787 mm central $\beta_{\perp}$

→ Excellent tool to systematically test novel density methods

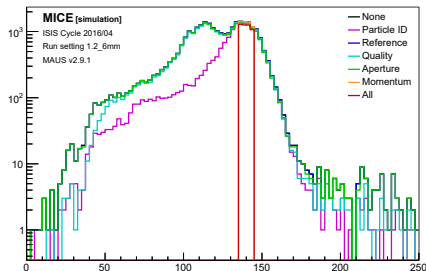
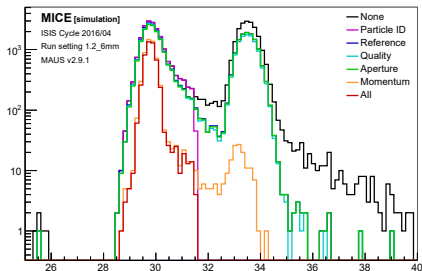


Particle selection

Series of cuts applied to both data and simulation:

- Muon tagging using TOF01 (**Particle ID**)
- Upstream reference plane hit (**Reference**)
- Good track reconstruction quality (**Quality**)
- Track within the tracker fiducial (**Aperture**)
- Long. momentum $\in [135, 145]$ MeV/c (**Momentum**)

All



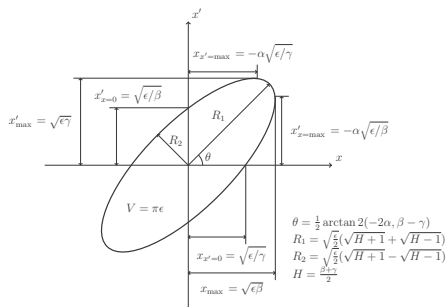
Transverse normalised RMS emittance

4D normalised RMS emittance:

$$\epsilon_n = \frac{1}{m} |\Sigma|^{\frac{1}{4}}, \quad (1)$$

with D the determinant of the covariance matrix defined as

$$\Sigma = \begin{pmatrix} \sigma_{xx} & \sigma_{xp_x} & \sigma_{xy} & \sigma_{xp_y} \\ \sigma_{p_x x} & \sigma_{p_x p_x} & \sigma_{p_x y} & \sigma_{p_x p_y} \\ \sigma_{yx} & \sigma_{yp_x} & \sigma_{yy} & \sigma_{yp_y} \\ \sigma_{p_y x} & \sigma_{p_y p_x} & \sigma_{p_y y} & \sigma_{p_y p_y} \end{pmatrix}. \quad (2)$$

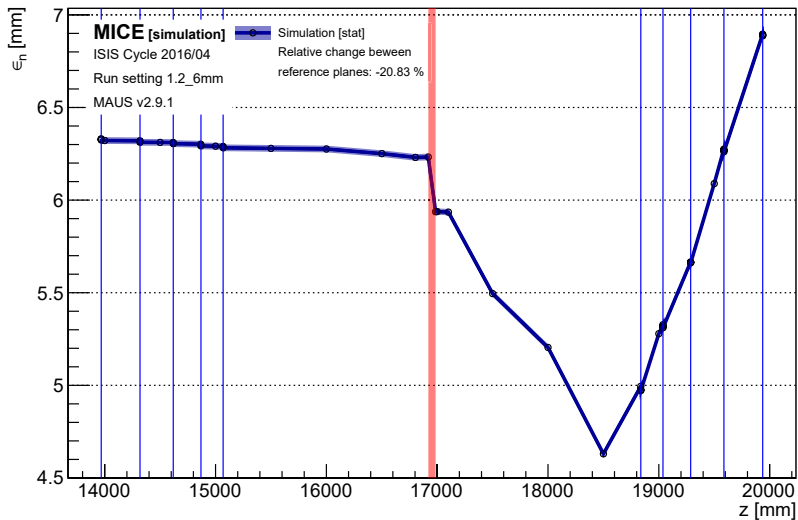


The RMS emittance is directly related to the **volume of the RMS ellipsoid** through $\epsilon_n = \sqrt{2V_{\text{RMS}}}/(m\pi)$ and as such is the most common probe of average phase-space density:

$$\rho_{\text{RMS}} = \frac{N}{V_{\text{RMS}}} = \frac{N}{\frac{1}{2}m^2\pi^2\epsilon_n^2} = \frac{N}{\frac{1}{2}\pi^2|\Sigma|^{\frac{1}{2}}} \quad [\text{mm}^{-2}(\text{MeV}/c)^{-2}]. \quad (3)$$

→ It follows from Liouville's theorem that the phase-space volume should be conserved

Emittance evolution



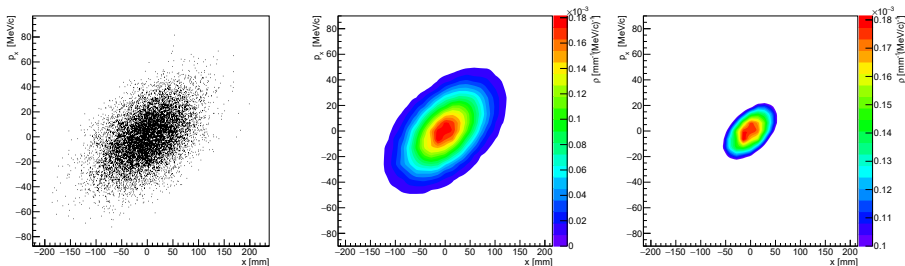
Power of density estimation

The emittance plot exhibits two obvious **challenges**:

- transmission losses yield apparent emittance reduction;
- filamentation in SSD yield apparent emittance growth.

The key to solving both problems lies in **density estimation**:

- Estimate density in the transverse 4D phase-space (*center*);
 - Select an **identical fraction** of the beam upstream and downstream from within the **densest area of the space** (*right*);
 - Define cooling figure of merits on these subsamples;
- The core, unlike the tails, is transmitted and linear.



Transverse single-particle amplitude

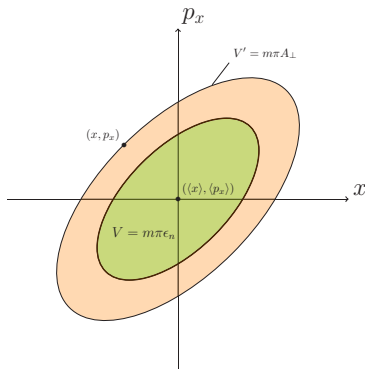
Single particle amplitude is defined as

$$A_{\perp} = \epsilon_n \mathbf{u}^T \Sigma^{-1} \mathbf{u} \quad (4)$$

with $\mathbf{u} = \mathbf{v} - \boldsymbol{\mu}$, the centered phase-space vector ($\mathbf{v}$) of the particle.

Amplitude follows a χ^2 **distribution** with d degrees of freedom with $\langle A_{\perp} \rangle = 4\epsilon_n$.

It is related to the **volume** of an ellipse similar to the RMS ellipse, going through $\mathbf{v}$.



Particle amplitude provides a density estimate in every input point

$$\rho(\mathbf{v}_i) = \frac{1}{(2\pi)^2 |\Sigma|^{\frac{1}{2}}} \exp \left[-\mathbf{u}^T \Sigma^{-1} \mathbf{u} / 2 \right] = \boxed{\frac{1}{4\pi^2 m^2 \epsilon_n^2} \exp \left[-\frac{A_{\perp}}{2\epsilon_n} \right]}. \quad (5)$$

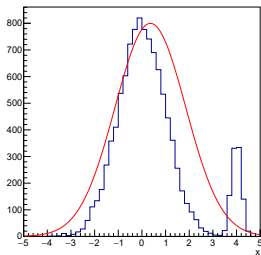
→ Allows for the selection of a **high density core** !

Amplitude reconstruction

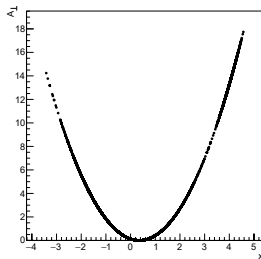
In the case of non-linear beams, special care must be taken in the reconstruction of amplitude as tails significantly bias the covariance matrix

Optimal procedure for amplitude reconstruction:

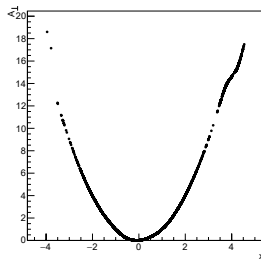
- Compute Σ and μ for the whole sample;
- 1 Calculate all the particle amplitudes $A_{\perp}^i$;
- 2 Register the highest amplitude in the distribution;
- 3 Update Σ and μ by removing the highest amplitude point;
- 4 Iterate from 1.



Test Gaussian + outliers

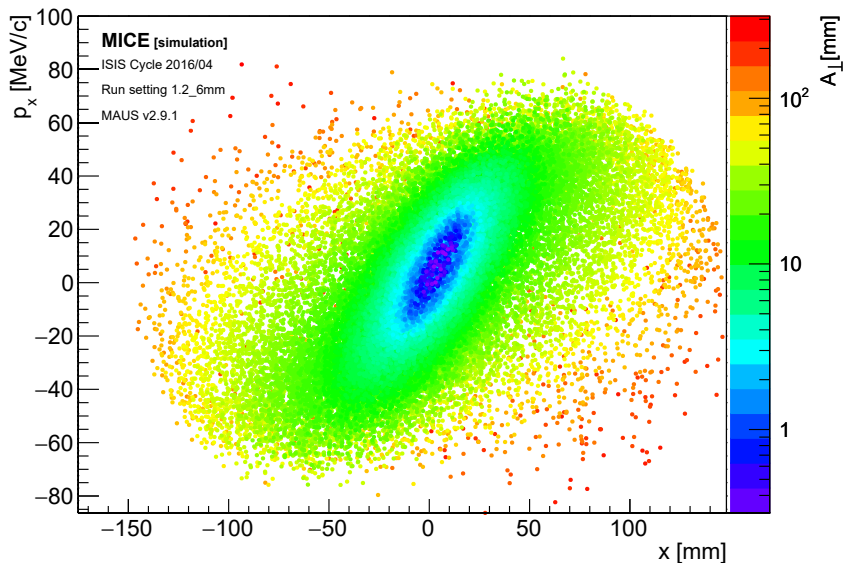
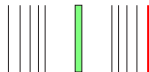


Regular amplitudes (biased)

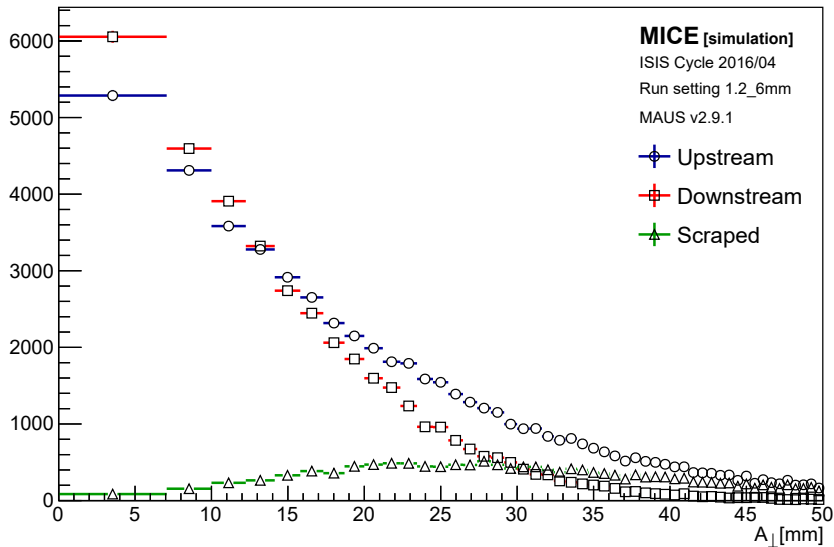
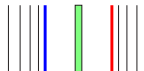


Corrected amplitudes

Amplitudes at TKD station 5



Amplitude distribution evolution

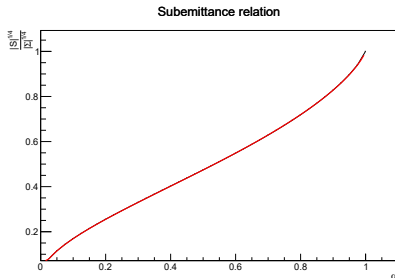


Subemittance definition and properties

The α -**subemittance**, e_α , is defined as the emittance of the core fraction α of the parent beam. For a truncated 4D Gaussian beam of covariance matrix S , it satisfies

$$\frac{e_\alpha}{\epsilon_n} = \frac{|S|^{\frac{1}{4}}}{|\Sigma|^{\frac{1}{4}}} = \frac{1}{2\alpha} \gamma\left(3, \frac{R^2}{2}\right), \quad (6)$$
$$R^2 = Q_{\chi_4^2}(\alpha).$$

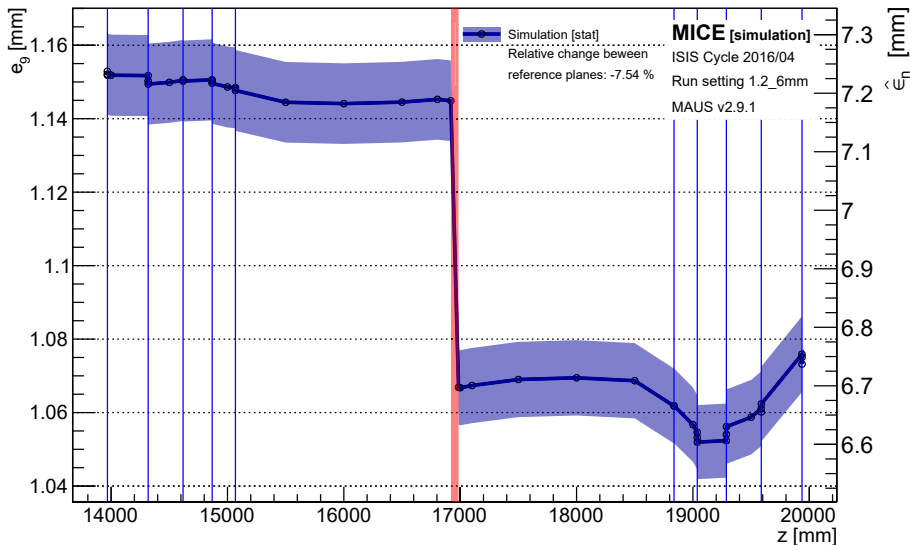
$$\rightarrow \boxed{\frac{e_\alpha^{\text{out}} - e_\alpha^{\text{in}}}{e_\alpha^{\text{in}}} = \frac{\epsilon_n^{\text{out}} - \epsilon_n^{\text{in}}}{\epsilon_n^{\text{in}}}}. \quad (7)$$



The statistical uncertainty carried by this measurement is identical to that of the emittance, scaled by the fraction α as

$$\frac{\sigma_{e_\alpha}}{e_\alpha} = \frac{1}{\sqrt{\alpha}} \frac{\sigma_{\epsilon_n}}{\epsilon_n} = \sqrt{\frac{2}{\alpha N d}}. \quad (8)$$

Submittance evolution



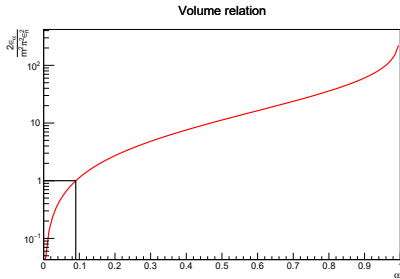
Fractional emittance definition and properties

The α -**fractional emittance**, ϵ_α , is defined as the phase-space volume occupied by the core fraction α of the parent beam. For a truncated 4D Gaussian beam of covariance matrix S , it satisfies

$$\epsilon_\alpha = \frac{1}{2} m^2 \pi^2 \epsilon_n^2 R^4 = V_{\text{RMS}} R^4, \quad (9)$$

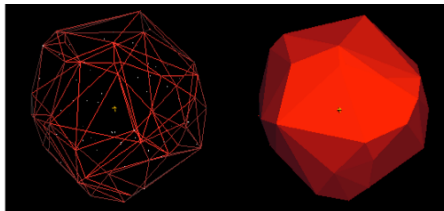
$$R^2 = Q_{\chi_4^2}(\alpha).$$

$$\rightarrow \boxed{\frac{\epsilon_\alpha^{\text{out}} - \epsilon_\alpha^{\text{in}}}{\epsilon_\alpha^{\text{in}}} \simeq 2 \frac{\epsilon_n^{\text{out}} - \epsilon_n^{\text{in}}}{\epsilon_n^{\text{in}}}} \quad (10)$$

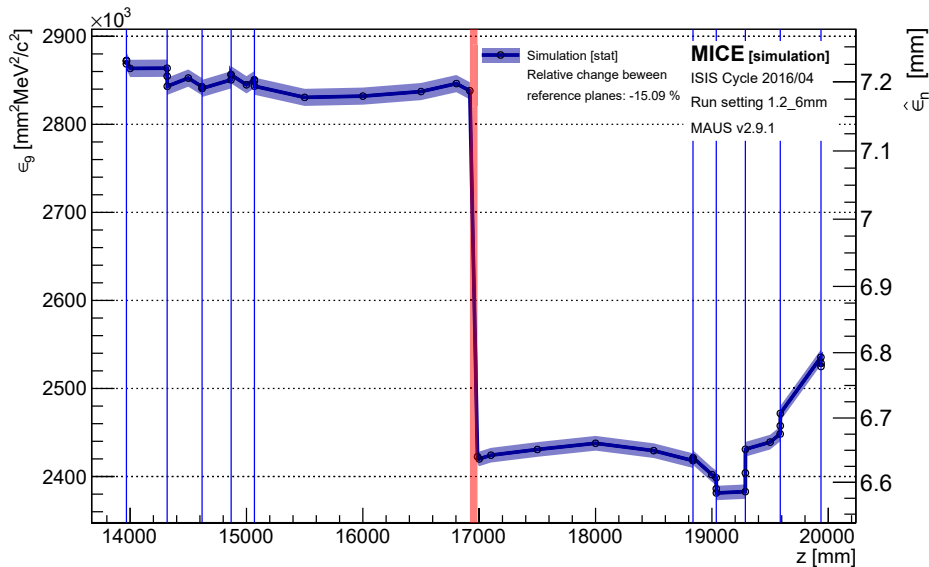


In 4D, a fraction α of **9%** yields the volume of the **RMS ellipsoid**, V_{RMS}

The **convex hull** is a prime candidate for volume reconstruction. It computes the smallest volume that contains the core αN points.



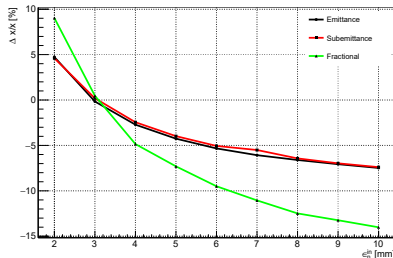
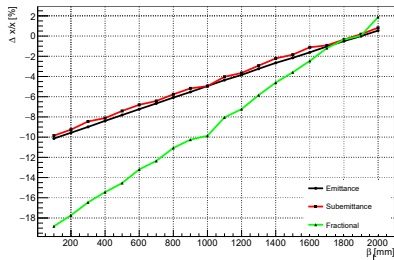
Phase-space volume evolution



Toy analysis of fractional quantities

A toy analysis (Gaussian input beam, toy absorber) shows:

- The **same relative change** is seen in the RMS emittance and all of the fractional quantities, for any fraction
- The change in fractional quantities exhibit the **same relation** with $\beta_{\perp}$ and the input emittance, ϵ_i
- The fractional quantities are **more robust** against losses and non-linearities as the tails do not influence their measurement



→ Plots produces for a core 9% selection, i.e. size of the RMS ellipse

Non-parametric density estimation: k -Nearest Neighbours

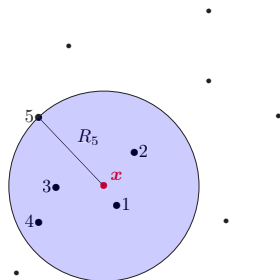
For a given point x , find the k **closest points** in the input cloud. Find the distance R_k to the k^{th} point and compute the 4D local density estimate as

$$\rho(x) = \frac{k}{\mathcal{V}_k} = \frac{2k}{\pi^2 R_k^4}, \quad (11)$$

with $\mathcal{V}_k$ the volume of the 4-ball centred in x of radius R_k .

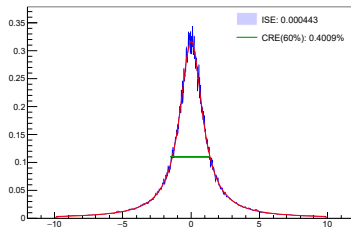
The rule of thumb choice of $k = \sqrt{N}$ yields quasi-optimal results for a broad array of distributions.

Right plot shows great agreement between theoretical Cauchy distribution (red) and estimation (blue).

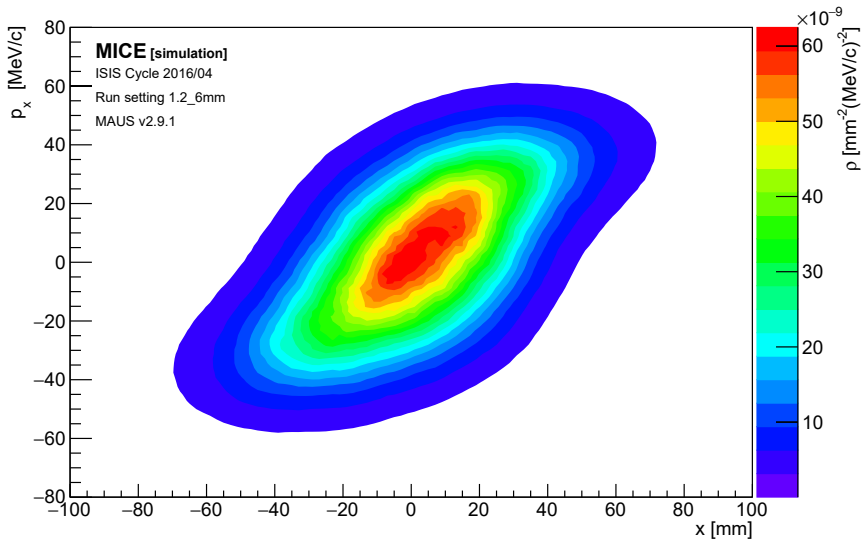
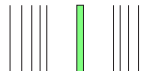


$$\hat{\rho}(x) = \frac{k}{V_k} = \frac{5}{\pi R_5^2}$$

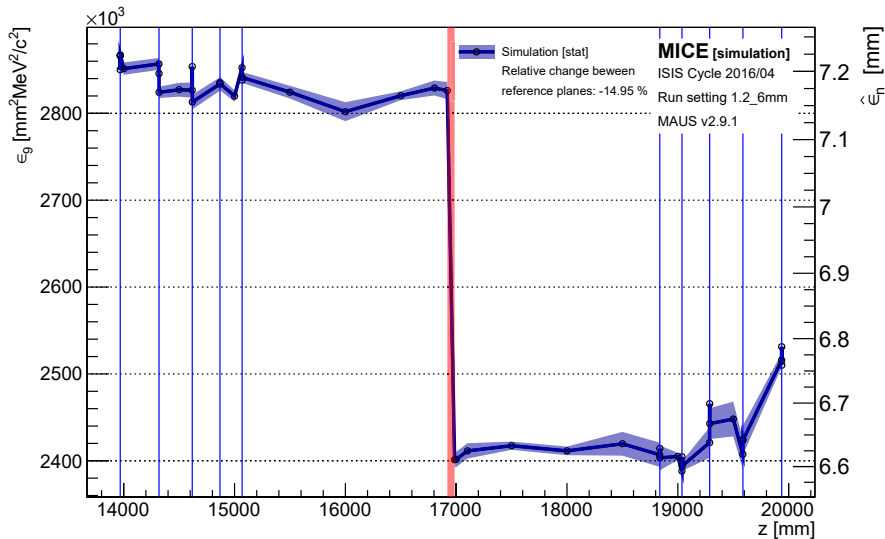
1D Cauchy distribution



Density estimation at TKD station 5



Phase-space volume evolution



Conclusions

Status of the amplitude-based analysis:

- Selecting the low amplitude core **gets rid of apparent emittance reduction** due to scraping and **apparent emittance growth** due to beam filamentation in the downstream section;
- A toy MC shows that the **exact same behaviour** can be observed for the subsample and fractional emittance as for the RMS definition;
- Method shows a **clean cooling signal** in a realistic MC.

Status of the non-parametric analysis:

- Systematic study well advanced, **k NN robust in 4D**, low error and no bias for large samples with the rule-of-thumb k selection;
- Method applied to the toy MC to study its behaviour, **identical trend** as with the amplitude-based fractional emittance;
- Method also shows **cooling signal** in a realistic MC.